

THE ARGMIN PROCESS OF BROWNIAN MOTION AND BROWNIAN EXTREMA OF A GIVEN LENGTH

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ABSTRACT. In this paper we investigate the argmin process of Brownian motion B defined by $\alpha_t := \sup \{s \in [0, 1] : B_{t+s} = \min_{u \in [0, 1]} B_{t+u}\}$ for $t \geq 0$. The argmin process α is stationary, with invariant measure which is arcsine distributed. We prove that $(\alpha_t; t \geq 0)$ is a Markov process with the Feller property, and provide its transition kernel $Q_t(x, \cdot)$ for $t > 0$ and $x \in [0, 1]$. In the second part of the paper we consider Brownian extrema of a given length. We prove that these extrema form a delayed renewal process with an explicit path construction. We also give a path decomposition for Brownian motion at these extrema.

Key words : arcsine law, argmin process, Brownian extrema, Feller semigroup, Brownian excursion theory, jump process, Lévy system, Markov property, space-time shift process, path decomposition, renewal property, sample path property, stationary process.

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1. INTRODUCTION AND MAIN RESULTS

In this paper we are interested in the *argmin process* $(\alpha_t; t \geq 0)$ of standard Brownian motion $(B_t; t \geq 0)$. That is,

$$\alpha_t := \sup \left\{ s \in [0, 1] : B_{t+s} = \min_{u \in [0, 1]} B_{t+u} \right\} \quad \text{for all } t \geq 0. \quad (1.1)$$

For each $t \geq 0$, $t + \alpha_t$ is the last time at which the minimum of B on $[t, t + 1]$ is achieved. The argmin process α is càdlàg, and takes values in $[0, 1]$. Except possible upward jumps, the process $(\alpha_t; t \geq 0)$ drifts down at unit speed. So $(\alpha_t; t \geq 0)$ can be interpreted as a storage process, see Çinlar and Pinsky [9, 10]. See also Çinlar [11], Brockwell et al. [7], and

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Evans and Pitman [18] for other examples of storage processes. The argmin process α also appeared as the hydrodynamic limit of a surface growth model in Belitsky and Ferrari [2].

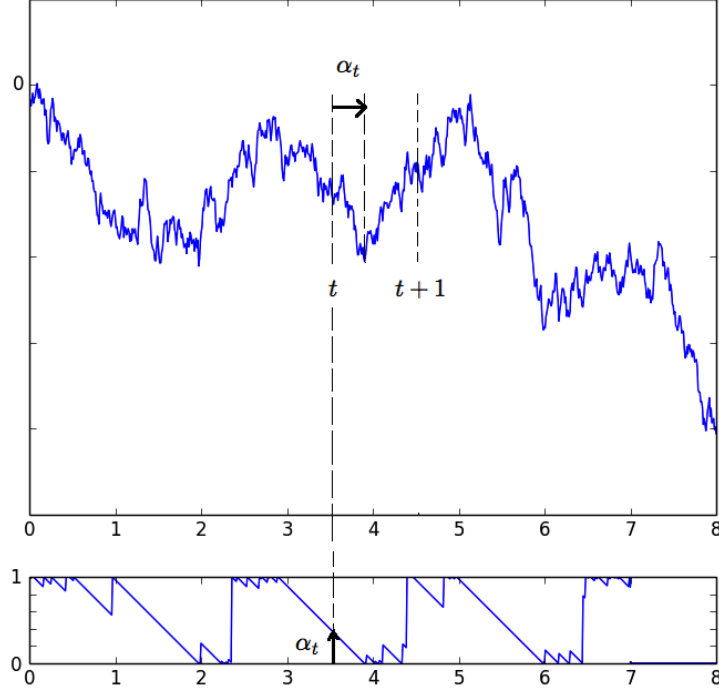


FIGURE 1. TOP: The argmin process embedded in Brownian motion. BOTTOM: The argmin process corresponding to the Brownian path on the top.

Let $(\Theta_t; t \geq 0)$ be the $\mathcal{C}[0, \infty)$ -valued space-time shift of Brownian motion B , defined by

$$\Theta_t := (B_{t+u} - B_t; u \geq 0) \quad \text{for all } t \geq 0.$$

In the path space setting,

$$\alpha_t = \alpha_0 \circ \Theta_t \quad \text{for all } t \geq 0. \quad (1.2)$$

By the stationarity of $(\Theta_t; t \geq 0)$, for any measurable function $g : \mathcal{C}[0, \infty) \rightarrow \mathbb{R}$, the process $(g \circ \Theta_t; t \geq 0)$ is stationary. Moreover, it is a well-known result of Lévy [33] that the time at which the minimum of Brownian motion on $[0, 1]$ is achieved follows the arcsine distribution. As a result, we have the following proposition.

Proposition 1.1. *The argmin process $(\alpha_t; t \geq 0)$ is stationary. The invariant measure is the arcsine distribution with density*

$$f(x) = \frac{1}{\pi \sqrt{x(1-x)}} \quad \text{for } 0 < x < 1. \quad (1.3)$$

It is natural to ask whether $(\alpha_t; t \geq 0)$ is a Markov process. Sufficient conditions for a function of a Markov process to be Markov are given by Dynkin [16] and Rogers and Pitman [46]. But these criteria do not apply to the argmin process, see Section 3.6. Nevertheless, we prove the following theorem.

Theorem 1.2. *The argmin process $(\alpha_t; t \geq 0)$ is a Markov process with Feller transition semigroup $Q_t(x, \cdot)$, $t > 0$ and $x \in [0, 1]$ where*

$$Q_t(x, dy) = \begin{cases} \frac{1_{\{0 < y < 1\}}}{\pi \sqrt{y(1-y)}} dy & \text{for } 0 \leq x \leq 1 < t, \\ \sqrt{\frac{1-x}{1-x+t}} \delta_{x-t}(dy) + \frac{\sqrt{(y+t-1)^+}}{\pi(y+t-x)\sqrt{1-y}} dy & \text{for } 0 < t \leq x \leq 1 \\ \frac{\sqrt{(1-x)(t-x)} + \sqrt{y(y+t-1)^+}}{\pi(y+t-x)\sqrt{y(1-y)}} dy & \text{for } 0 \leq x < t \leq 1. \end{cases} \quad (1.4)$$

See Kallenberg [27, Chapter 19] for background on Feller semigroups of continuous-time Markov processes. The proof of Theorem 1.2 is given in Sections 3.2 and 3.4.

Our approach relies on Brownian excursion theory, see Section 2. We also investigate the law of jumps of $(\alpha_t; t \geq 0)$. In particular, we prove that the argmin process α has local times at levels 0 and 1, and provide a Lévy system of $(\alpha_t; t \geq 0)$. As a consequence of Theorem 1.2, we compute the infinitesimal generator of the argmin process α , which leads to a martingale characterization of $(\alpha_t; t \geq 0)$. These results imply that the argmin process α is a time-homogeneous Markov process with an explicit description in the framework of *jumping Markov processes* by Jacod and Skorokhod [26], following the study of *piecewise deterministic Markov processes* by Davis [13].

It will be seen that the argmin process of Brownian motion has quite an intricate structure, which we prefer to analyze in detail before discussing various possible generalizations. In a forthcoming article, we will

- use Millar's decomposition [36] to show that the argmin process of a Lévy process is Markov;
- calculate explicitly the semigroup of the argmin process of a stable process;
- study a discrete analog of the argmin process: the argmin chain of a random walk.

Our motivation for considering the argmin process comes from the study of Brownian extrema with given length. In the second part of this paper we provide further insight into these extrema, following previous works of Neveu and Pitman [38], and Leuridan [32]. For $a, b > 0$, let

$$\mathcal{M}_{a,b} := \left\{ t \geq a : B_t = \min_{s \in [t-a, t+b]} B_s \right\} = \{T_1^{a,b}, T_2^{a,b}, \dots\}, \quad (1.5)$$

with $a < T_1^{a,b} < T_2^{a,b} < \dots$ be the (a, b) -minima set of Brownian motion B .

The study of Brownian extrema dates back to Lévy [33]. See Freedman [20, Section 1.4], Karatzas and Shreve [28, Section 2.9] for development. Neveu and Pitman [38] proved the renewal property of Brownian local extrema by looking at Brownian extrema of a given depth. They gave the following Palm description of Brownian local extrema.

Theorem 1.3. [38] *Let $(\mathcal{C}, \mathcal{B})$ be the space of continuous paths on \mathbb{R} , equipped with Wiener measure \mathbf{W} , and (E, \mathcal{E}) be the space of excursions with lifetime ζ , equipped with Itô's law \mathbf{n} (see Section 2.2 for discussion). For $A \in \mathcal{B}$, define the σ -finite Palm measure of Brownian local minima by*

$$\nu(A) := \mathbb{E} \# \{0 \leq t \leq 1 : t \text{ is a local minimum and } \theta_t \in A\}, \quad (1.6)$$

where $\theta_t := (b_{t+u} - b_t; t \in \mathbb{R})$ is the spacetime shift of a two-sided Brownian motion b . Then for $A \in \mathcal{B}$,

$$\nu(A) = \frac{1}{2}(\mathbf{n} \otimes \mathbf{n} \otimes \mathbf{W})(f^{-1}(A)), \quad (1.7)$$

where $f : E \times E \times \mathcal{C} \ni (e, e', w) \rightarrow \tilde{w} \in \mathcal{C}$ is a mapping given by

$$\tilde{w}_t = \begin{cases} w_{t+\zeta(e')} & \text{if } t \leq -\zeta(e'), \\ e'_{-t} & \text{if } -\zeta(e') \leq t \leq 0, \\ e_t & \text{if } 0 \leq t \leq \zeta(e), \\ w_{t-\zeta(e)} & \text{if } t \geq \zeta(e). \end{cases} \quad (1.8)$$

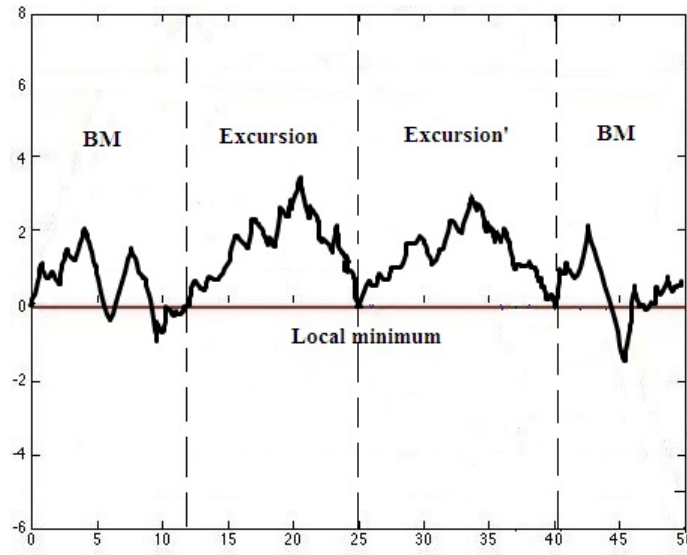


FIGURE 2. Structure of a Brownian local minimum.

The quantity $\nu(A)$ is interpreted as the mean number of Brownian local minima of type A per unit time. See Kallenberg [27, Chapter 11] for background on Palm measures.

Neveu-Pitman's results were generalized to Brownian motion with drift by Faggionato [19]. Inspired from [38], Leuridan [32] considered Brownian extrema with given length. In different directions, Groeneboom [22] considered the global extremum of Brownian motion with a parabolic drift, where he gave a density formula in terms of *Airy functions*. Tsirelson [49] provided an i.i.d. uniform sampling construction of Brownian local extrema under external randomization. Abramson and Evans [1] considered Lipschitz minorants of Brownian motion, which is a variant of Brownian extrema.

Leuridan studied the (a, b) -minima set $\mathcal{M}'_{a,b}$ of a two-sided Brownian motion. He proved that times of the set $\mathcal{M}'_{a,b}$ form a renewal process, and provided the density of the inter-arrival times. Observe that

$$\mathcal{M}_{a,b} - a \stackrel{(d)}{=} \mathcal{M}'_{a,b} \cap [0, \infty). \quad (1.9)$$

That is, $\mathcal{M}_{a,b} - a := \{T_i^{a,b} - a; i \geq 1\}$ is a renewal process with stationary delay. We adapt Leuridan's result to one-sided Brownian motion as follows.

Theorem 1.4. [32] *Let $a, b > 0$. The times of (a, b) -minima of Brownian motion $(B_t; t \geq 0)$ form a delayed renewal process, denoted by $(T_i^{a,b}; i \geq 1)$ so that $a < T_1^{a,b} < T_2^{a,b} < \dots$. Let*

$$h_{a,b}(t) := \begin{cases} \frac{1}{\pi t} \left(\sqrt{\frac{(t-b)^+}{b}} + \sqrt{\frac{(t-a)^+}{a}} \right) & \text{for } 0 < t < a+b \\ \frac{1}{\pi\sqrt{ab}} & \text{for } t \geq a+b. \end{cases} \quad (1.10)$$

Then $(\Delta_i^{a,b} := T_{i+1}^{a,b} - T_i^{a,b}; i \geq 1)$ are independent, with density

$$g_{a,b}(t) := \sum_{n=1}^{\infty} (-1)^{n-1} h_{a,b}^{*n}(t), \quad (1.11)$$

where $h_{a,b}^{*n}$ is the n^{th} convolution of $h_{a,b}$. In addition, $T_1^{a,b}$ is independent of $(\Delta_i^{a,b}; i \geq 1)$, and has density

$$f_{a,b}(t) := \frac{1_{\{t>a\}}}{\pi\sqrt{ab}} \int_{t-a}^{\infty} g_{a,b}(s) ds. \quad (1.12)$$

Given a measurable set $A \subset \mathbb{R}^+$, let $N_{a,b}(A) := \#(\mathcal{M}_{a,b} \cap A)$ be the counting measure of (a, b) -minima in Brownian motion B . Leuridan's proof of Theorem 1.4 is based on the formula, for $n \geq 1$ and $0 < t_1 < \dots < t_n$,

$$\mathbb{E}(N_{a,b}(dt_1) \cdots N_{a,b}(dt_n)) = \frac{1}{\pi\sqrt{ab}} \prod_{k=1}^{n-1} h_{a,b}(t_k - t_{k-1}) dt_1 \cdots dt_n, \quad (1.13)$$

with convention that $\prod_{\emptyset} := 1$. The case $n = 1$ of (1.13) follows readily from Theorem 1.3, since for a generic (a, b) -minimum, the left excursion has length larger than a and the right excursion has length larger than b . This implies that the mean number of (a, b) -minima per unit time is given by

$$\frac{1}{2} \mathbf{n}(\zeta(e') > a) \mathbf{n}(\zeta(e) > b) = \frac{1}{\pi\sqrt{ab}}.$$

In particular,

$$\mathbb{E}(\Delta_i^{a,b}) = \pi\sqrt{ab} \quad \text{for all } i \geq 1. \quad (1.14)$$

However, to obtain (1.13) for $n \geq 2$ requires extra work. Observe that for $a + b = 1$, the set $\mathcal{M}_{a,b}$ can be viewed as the a -level set of the argmin process α . So by Brownian scaling,

$$\begin{aligned} \mathcal{M}_{a,b} &\stackrel{(d)}{=} (a+b) \mathcal{M}_{\frac{a}{a+b}, \frac{b}{a+b}} \\ &\stackrel{(d)}{=} (a+b) \alpha^{-1} \left(\left\{ \frac{a}{a+b} \right\} \right) \quad \text{for } a, b > 0. \end{aligned} \quad (1.15)$$

According to Hoffmann-Jørgensen [23], and Krylov and Juškevič [31], the set $\mathcal{M}_{a,b}$ enjoys regenerative property, and is called a *strong Markov set*. See also Kingman [29] for a survey on regenerative phenomena of level sets of Markov processes. In Section 4.1, we recover Theorem 1.4, in particular (1.13), by using the properties of the argmin process α .

Note that the density $h_{a,b}$ defined by (1.10) is induced by a σ -finite measure. By Leuridan's formula (1.11), the Laplace transform of $\Delta_i^{a,b}$ is given by

$$\Phi_{\Delta_{a,b}}(\lambda) = \frac{\Psi(\lambda)}{1 + \Psi(\lambda)} \quad \text{with } \Psi(\lambda) := \int_0^\infty e^{-\lambda t} h_{a,b}(t) dt, \quad (1.16)$$

provided that $\Psi(\lambda) < 1$. By analytic continuation, we extend (1.16) to all $\lambda > 0$. But it does not seem obvious how to simplify (1.16) analytically.

While the description of $\mathcal{M}_{a,b}$ is complicated for general $a, b > 0$, the case $a = b$ is simplified. For simplicity, we consider $a = b = 1$. We shall give a construction of the $(1, 1)$ -minima set

$$\mathcal{M}_{1,1} := \{T_1, T_2, \dots\} \quad \text{with } 1 < T_1 < T_2 < \dots,$$

from which we derive simple formulas for the Laplace transforms of $\Delta_i := T_{i+1} - T_i$ and T_1 .

Let J be the first descending ladder time of Brownian motion, from which starts an excursion above the minimum of length larger than 1. It is known that the Laplace transform of J is given by

$$\Phi_J(\lambda) = \frac{1}{\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda}}, \quad (1.17)$$

where $\operatorname{erf}(x) := \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$ is the error function. See Proposition 2.3 for basic properties of J .

The random variable J plays an important role in our construction of the $(1, 1)$ -minima set. Let Δ be distributed as the law of the inter-arrival times Δ_i , independent of J . It is a simple consequence of the construction in Section 4.2 of the Brownian path over $[0, T_1]$ and over $[T_i, T_{i+1}]$ that

$$T_1 \stackrel{(d)}{=} J + 1_{\{J < 1\}} \Delta, \quad \text{with } J \text{ independent of } \Delta. \quad (1.18)$$

Combined with the fact that $T_1 - 1$ is the stationary delay for a renewal process with inter-arrival time distributed according to Δ , this leads to the following result:

Theorem 1.5. *Let $(T_i; i \geq 1)$ with $1 < T_1 < T_2 < \dots$ be times of the $(1, 1)$ -minima set $\mathcal{M}_{1,1}$ of Brownian motion B .*

- (1) *Let J' be an independent copy of J , whose Laplace transform is given by (1.17). Then there is the identity in law*

$$T_1 - 1 \stackrel{(d)}{=} J + J'. \quad (1.19)$$

In particular, the Laplace transforms of T_1 and Δ are given by

$$\Phi_{T_1}(\lambda) = e^{-\lambda} (\Phi_J(\lambda))^2 = \frac{e^{-\lambda}}{(\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda})^2}, \quad (1.20)$$

and

$$\Phi_\Delta(\lambda) = 1 - \pi\lambda (\Phi_J(\lambda))^2 = 1 - \frac{\pi\lambda}{(\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda})^2}. \quad (1.21)$$

Moreover,

$$\mathbb{E}T_1 = 3 \quad \text{and} \quad \mathbb{E}\Delta = \pi. \quad (1.22)$$

(2) The fragments $(B_{T_i+t} - B_{T_i}; 0 \leq t \leq \Delta_i)_{i \geq 1}$, are i.i.d., starting as Brownian meander of length 1, and then running as Brownian motion until the next $(1, 1)$ -minima occurs. Moreover,

$$\mathbb{E}B_{T_1} = -\sqrt{\frac{\pi}{2}} \quad \text{and} \quad \text{Var } B_{T_1} = \frac{4 + \pi}{2}, \quad (1.23)$$

and

$$\mathbb{E}(B_{T_{i+1}} - B_{T_i}) = 0 \quad \text{and} \quad \text{Var}(B_{T_{i+1}} - B_{T_i}) = \pi \quad \text{for all } i \geq 1. \quad (1.24)$$

In Section 4.2, we prove the identity in law (1.19) by computing the Laplace transform (1.20) of T_1 . This identity in law is surprising, and we do not have a simple explanation. Though we are able to compute the first two moments (1.23)-(1.24), the laws of B_{T_1} , and $B_{T_{i+1}} - B_{T_i}$ seem to be difficult. We leave these for further investigation.

Let $\mathcal{M}'_{1,1} := \{\dots, T'_0, T'_1, \dots\}$ with $\dots < T'_0 < 0 \leq T'_1$ be the $(1, 1)$ -minima set of a two-sided Brownian motion $(b_t; t \in \mathbb{R})$. By (1.9), the identity (1.19) is equivalent to

$$T'_1 \stackrel{(d)}{=} J + J'.$$

Note that

$$B_{T_1} = B_1 + (B_{T_1} - B_1) \quad \text{with } B_{T_1} - B_1 \stackrel{(d)}{=} b_{T'_1},$$

and

$$B_{T_{i+1}} - B_{T_i} \stackrel{(d)}{=} b_{T'_{i+1}} - b_{T'_i} \quad \text{for all } i \geq 1.$$

But B_1 and $B_{T_1} - B_1$ are not independent. Then (1.23) and (1.24) imply that

$$\mathbb{E}b_{T'_1} = -\sqrt{\frac{\pi}{2}}, \quad (1.25)$$

and

$$\mathbb{E}(b_{T'_{i+1}} - b_{T'_i}) = 0 \quad \text{and} \quad \text{Var}(b_{T'_{i+1}} - b_{T'_i}) = \pi \quad \text{for all } i \geq 1. \quad (1.26)$$

Recall the notations in Theorem 1.3. For $a > 0$, let

$$\nu^a(A) := \mathbb{E}\#\{0 \leq t \leq 1 : t \in \mathcal{M}'_{a,a} \text{ and } \theta_t \in A\} \quad \text{for } A \in \mathcal{B},$$

be the Palm measure of (a, a) -minima of a two-sided Brownian motion. Theorem 1.5 implies that ν^a has total mass $\frac{1}{\pi a}$, and

$$\nu^a(A) = \frac{1}{2} 1_{\{\zeta' > a, \zeta > a\}}(\mathbf{n} \otimes \mathbf{n} \otimes \mathbf{W})(f^{-1}(A)) \quad \text{for } A \in \mathcal{B}, \quad (1.27)$$

where f is the mapping defined by (1.8). By taking $a \downarrow 0$, the Palm measures ν^a increase to the limit ν defined by (1.7). This recovers Theorem 1.3.

Finally, the set $\mathcal{M}_{1,1}$ is directly related to the argmin process α without scaling. In fact, T_i is the i^{th} time that the process α reaches 0 by a continuous passage from 1. So the law of Brownian fragments between $(1, 1)$ -minima can be derived from the study of α . Let

$$\mathcal{LE} := \{t \geq 0 : B_t < B_s, \text{ for all } s \in [t, t+1]\},$$

be the left ends of forward meanders of length 1, and

$$\mathcal{RE} := \{t \geq 1 : B_t < B_s, \text{ for all } s \in [t-1, t]\},$$

be the right ends of backward meanders of length 1. In Lemma 4.5, we show that left ends come before right ends between any two consecutive $(1, 1)$ -minima. So we define for each $i \geq 1$,

$$D_i := \inf\{t > T_i : t \in \mathcal{RE}\} \quad \text{and} \quad G_i := \sup\{t \in D_i : t \in \mathcal{LE}\}. \quad (1.28)$$

For each $i \geq 1$, the triple $(G_i - T_i, D_i - G_i, T_{i+1} - D_i)$ gives a decomposition of Δ_i :

$$\Delta_i = (G_i - T_i) + (D_i - G_i) + (T_{i+1} - D_i). \quad (1.29)$$

By using the Lévy system of the argmin process, we prove the following theorem which identifies the law of this triple.

Theorem 1.6. *For each $i \geq 1$, $G_i - T_i$, $D_i - G_i$ and $T_{i+1} - D_i$ are mutually independent, with*

- $G_i - T_i \stackrel{(d)}{=} T_{i+1} - D_i \stackrel{(d)}{=} J$, the Laplace transform of which is given by (1.17);
- the density of $D_i - G_i$ is given by

$$\mathbb{P}(D_i - G_i \in dt) = \frac{2-t}{t^2\sqrt{t-1}} 1_{\{1 < t < 2\}} dt. \quad (1.30)$$

In particular, for each $i \geq 1$,

$$\mathbb{E}(G_i - T_i) = \mathbb{E}(T_{i+1} - D_i) = 1 \quad \text{and} \quad \mathbb{E}(D_i - G_i) = \pi - 2. \quad (1.31)$$

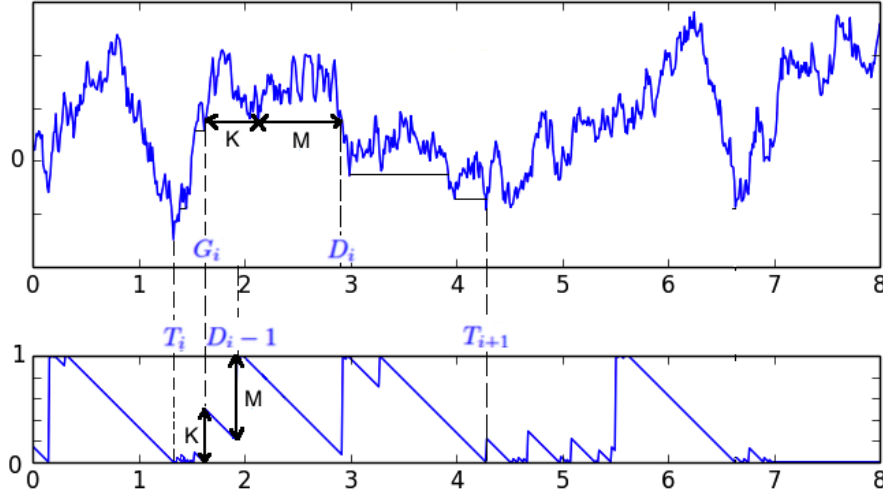


FIGURE 3. (T_i, G_i, D_i, T_{i+1}) related to the argmin process $(\alpha_t; t \geq 0)$.

For a random variable X , let $\Phi_X(\lambda)$ be the Laplace transform of X , and $f_X(\cdot)$ be the density of X . Theorems 1.4-1.6 provide three different descriptions of the inter-arrival time Δ . This leads to some non-trivial identities. We summarize the results in the following table.

	Laplace transform $\Phi_X(\lambda)$	Density $f_X(t)$
$X = J$, Prop 2.3	$\frac{1}{\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda}}$	Given by (2.11)
$X = D - G$, Th 1.6	$\pi\lambda \left[(\operatorname{erf}(\sqrt{\lambda})^2 - 1) \right] + 2\sqrt{\pi\lambda}e^{-\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-2\lambda}$	$\frac{2-t}{t^2\sqrt{t-1}} 1_{\{1 < t < 2\}}$
$X = \Delta$, Th 1.4	$\frac{\Psi(\lambda)}{1 + \Psi(\lambda)}$, with $\Psi(\lambda) = \operatorname{erf}(\sqrt{\lambda})^2 - 1 + \frac{2e^{-\lambda}}{\sqrt{\pi\lambda}} \operatorname{erf}(\sqrt{\lambda}) + \frac{e^{-2\lambda}}{\pi\lambda}$	Given by (1.11), with $a = b = 1$
$X = \Delta$, Th 1.5	$1 - \pi\lambda(\Phi_J(\lambda))^2$	$-\pi \frac{d}{dt} f_J^{*2}(t)$
$X = \Delta$, Th 1.6	$(\Phi_J(\lambda))^2 \Phi_{D-G}(\lambda)$	$(f_J * f_{D-G} * f_J)(t)$
$X = T_1 - 1$, Th 1.4	$\frac{1}{\pi\lambda} (1 - \Phi_\Delta(\lambda))$	$\frac{1_{\{t>0\}}}{\pi} \int_t^\infty f_\Delta(s) ds$
$X = T_1 - 1$, Th 1.5	$(\Phi_J(\lambda))^2$	$f_J^{*2}(t)$

TABLE 1. Laplace transforms and densities.

Organization of the paper: The layout of the paper is as follows.

- In Section 2, we provide background and necessary tools which will be used later.
- In Section 3, we study the argmin process $(\alpha_t; t \geq 0)$ of Brownian motion, and prove Theorem 1.2.
- In Section 4, we study the (a, b) -minima of Brownian motion with an emphasis on the case $a = b = 1$. There we prove Theorems 1.4, 1.5 and 1.6.

2. BACKGROUND AND TOOLS

This section recalls some background of Brownian motion and Markov processes. In Section 2.1, we consider Denisov's decomposition for Brownian motion. In Section 2.2, we recall various results from Brownian excursion theory.

2.1. Path decomposition of Brownian motion. Let $(B_t; t \geq 0)$ be standard Brownian motion. A Brownian meander $(m_t; 0 \leq t \leq 1)$ can be regarded as the weak limit of

$$\left(B_t; 0 \leq t \leq 1 \mid \inf_{0 \leq s \leq 1} B_s > -\epsilon \right) \quad \text{as } \epsilon \downarrow 0.$$

We refer to Durrett et al. [15] for a proof. A Brownian meander of length x , say $(m_t^x; 0 \leq t \leq x)$ is defined as

$$m_t^x := \sqrt{x} m_{t/x} \quad \text{for } 0 \leq t \leq x.$$

In particular, m_x^x is Rayleigh distributed with scale parameter \sqrt{x} . That is,

$$\mathbb{P}(m_x^x \in dy) = \frac{y}{x} \exp\left(-\frac{y^2}{2x}\right) dy \quad \text{for } y > 0. \quad (2.1)$$

In particular,

$$\mathbb{E} m_x^x = \sqrt{\frac{\pi x}{2}} \quad \text{and} \quad \text{Var } m_x^x = \frac{4 - \pi}{2} x. \quad (2.2)$$

The following path decomposition is due to Denisov.

Theorem 2.1 (Denisov's decomposition for Brownian motion). [14] *Let $A := \operatorname{argmin}_{u \in [0,1]} B_u$ be the time at which Brownian motion B attains its minimum on $[0, 1]$. Given A , which is arcsine distributed, the Brownian path is decomposed into two conditionally independent pieces:*

- (a). $(B_{A-t} - B_A; 0 \leq t \leq A)$ is a Brownian meander of length A ;
- (b). $(B_{A+t} - B_A; 0 \leq t \leq 1 - A)$ is a Brownian meander of length $1 - A$.

Let

$$\mathbf{P}^x := \overleftarrow{\mathbf{M}}^x \otimes \overrightarrow{\mathbf{M}}^{1-x} \otimes \mathbf{W} \quad \text{for } 0 \leq x \leq 1 \quad (2.3)$$

be the law of two independent Brownian meanders of length x and $1 - x$ joined back-to-back, concatenated by an independent Brownian path running forever. Denisov's decomposition is equivalent to

$$\mathbf{W}(\cdot) = \int_0^1 \frac{1}{\pi \sqrt{x(1-x)}} \mathbf{P}^x(\cdot) dx. \quad (2.4)$$

It is closely related to Williams' decomposition [51] by Brownian scaling, see Pitman and Ross [41, Section 2]. Denisov's decomposition can also be obtained by random walks approximation, see Iglehart [24] and Bolthausen [6].

2.2. Brownian excursion theory. Let $(B_t; t \geq 0)$ be standard Brownian motion, and $\underline{B}_t := \inf_{0 \leq s \leq t} B_s$ be the past-minimum process of B . For $l > 0$, let $T_l := \inf\{t > 0; B_t < -l\}$ be the first time at which B hits below level $-l$. Let

$$\mathcal{D} := \{l > 0 : T_{l-} < T_l\},$$

so that for $l \in \mathcal{D}$,

$$e_l := \begin{cases} B_{T_{l-}+t} - \underline{B}_{T_{l-}+t} & \text{for } 0 \leq t \leq T_l - T_{l-} \\ 0 & \text{for } t > T_l - T_{l-} \end{cases}$$

is an excursion away from $-l$. Let E be the space of excursions defined by

$$E := \{\epsilon \in \mathcal{C}[0, \infty); \epsilon_0 = 0, \epsilon_t > 0 \text{ for } t \in (0, \zeta(\epsilon)), \text{ and } \epsilon_t = 0 \text{ for } t \geq \zeta(\epsilon)\},$$

where $\zeta(\epsilon) := \inf\{t > 0; \epsilon_t = 0\} \in (0, \infty)$ is the lifetime of the excursion $\epsilon \in E$. The following theorem is a special case of Itô's excursion theory.

Theorem 2.2. [25] *The point measure*

$$\sum_{l \in \mathcal{D}} \delta_{(l, e_l)}(ds d\epsilon)$$

is a Poisson point process on $\mathbb{R}^+ \times E$ with intensity $ds \times \mathbf{n}(d\epsilon)$, where $\mathbf{n}(d\epsilon)$, called Itô's excursion law, is a σ -finite measure on E .

Here we consider positive excursions of the reflected process $B - \underline{B}$. So the measure $\mathbf{n}(d\epsilon)$ corresponds to $2\mathbf{n}^+(d\epsilon)$ defined in Revuz and Yor [45, Chapter VII].

Let $\Lambda(dx)$ be the Lévy measure of a $\frac{1}{2}$ -stable subordinator such that

$$\Lambda(dx) = \frac{dx}{\sqrt{2\pi x^3}} \quad \text{and} \quad \Lambda(x, \infty) = \sqrt{\frac{2}{\pi x}} \quad \text{for } x > 0. \quad (2.5)$$

By applying the master formula of Poisson point processes, we know that

$$\mathbf{n}(\zeta \in dx) = \Lambda(dx). \quad (2.6)$$

See Revuz and Yor [45, Chapter XII] for development of Brownian excursion theory. Let

$$A_t := \sum_l 1_{\{T_{l-} \leq t \leq T_l\}}(t - T_{l-}), \quad (2.7)$$

be the *age process* of excursions of $B - \underline{B}$, or equivalently of a $\frac{1}{2}$ -stable subordinator, and

$$R_t := \sum_l 1_{\{T_{l-} \leq t \leq T_l\}}(T_l - t), \quad (2.8)$$

be the *residual life* of excursions of $B - \underline{B}$, or equivalently of a $\frac{1}{2}$ -stable subordinator. Inspired by Krylov and Juškevič [31], Kingman [29] derived the joint distribution of (A_t, R_t) by the triple Laplace transform

$$\int_0^\infty e^{-\alpha t} \mathbb{E} e^{-\theta A_t - \phi R_t} = \frac{\sqrt{\phi} - \sqrt{\theta + \alpha}}{(\phi - \theta - \alpha)\sqrt{\alpha}} \quad \text{for } \phi \neq \theta + \alpha. \quad (2.9)$$

The following proposition gathers useful results of J , defined by

$$J + 1 := \inf\{t > 0 : A_t = 1\}.$$

That is, J is the first descending ladder time of Brownian motion, from which starts an excursion above the minimum of length exceeding 1. For completeness, we include a proof.

Proposition 2.3. [43] *Let J be the first descending ladder time of Brownian motion, from which starts an excursion above the minimum of length exceeding 1.*

- (1) *The random variable $\frac{1}{1+J}$ has the same distribution as the longest interval of Poisson-Dirichlet $(\frac{1}{2}, 0)$ distribution. The Laplace transform of J is given by (1.17), and*

$$\mathbb{E}J = 1. \quad (2.10)$$

- (2) *The density of J is given by*

$$\frac{\mathbb{P}(J \in dt)}{dt} = \sum_{n=1}^{\infty} (-1)^{n+1} \frac{c_n}{\sqrt{t}} I_n \left(\frac{1}{1+t} \right) \quad \text{for } t > 0, \quad (2.11)$$

where for each $n \geq 1$,

$$c_n := \frac{(n-1)!}{2^{n-1} \pi^{n/2} \Gamma(n/2)},$$

and I_n is a function supported on $(0, \frac{1}{n}]$ defined by

$$I_n(u_n) := \int \prod_{i=1}^{n-1} \sqrt{\frac{(1-u_i)^{n-1-i}}{u_i^3}} 1_{\left\{ \frac{u_{i+1}}{1-u_{i+1}} \leq u_i \leq \frac{1}{i} \right\}} du_i \quad \text{for } u_n \in \left(0, \frac{1}{n} \right],$$

with convention that $I_1(u_1) := 1$ for $u_1 \in (0, 1]$. In particular,

$$\mathbb{P}(J \in dt) = \frac{dt}{\pi \sqrt{t}} \quad \text{for } 0 < t \leq 1. \quad (2.12)$$

and

$$\mathbb{P}(J \in dt) = \frac{1}{\pi} \left(\frac{2}{\sqrt{t}} - 1 \right) dt \quad \text{for } 1 < t \leq 2, \quad (2.13)$$

Proof. The part (1) is essentially from Pitman and Yor [43, Corollary 12] with $\alpha = \frac{1}{2}$. Alternatively, let $\tau := \inf\{l \in D : T_l - T_{l-} > 1\}$ be the first level above which an excursion has length larger than 1 so that $J = T_{\tau-}$. As in [21], we deduce from Theorem 2.2 that τ is exponentially distributed with rate $\Lambda(1, \infty) = \sqrt{2/\pi}$, independent of $(B - \underline{B})[I_\tau]$ and that

$$J \stackrel{(d)}{=} \sigma_\xi,$$

where

- $(\sigma_t; t \geq 0)$ is a $\frac{1}{2}$ -stable subordinator with all jumps of size larger than 1 deleted, so the Laplace exponent of $(\sigma_t; t \geq 0)$ is given by

$$\phi(\lambda) := \int_0^1 (1 - e^{-\lambda x}) \Lambda(dx) = \sqrt{2\lambda} \operatorname{erf}(\sqrt{\lambda}) - \sqrt{\frac{2}{\pi}} (1 - e^{-\lambda}) \quad \text{for } \lambda \geq 0.$$

- ξ is exponentially distributed with rate $\sqrt{2/\pi}$, independent of $(\sigma_t; t \geq 0)$.

So

$$\mathbb{E}J = \mathbb{E}\sigma_1 \mathbb{E}\xi = \sqrt{\frac{2}{\pi}} \cdot \sqrt{\frac{\pi}{2}} = 1,$$

and the Laplace transform of J is given by

$$\Phi_J(\lambda) = \frac{\sqrt{2/\pi}}{\sqrt{2/\pi} + \phi(\lambda)} = \frac{1}{\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda}}.$$

The part (2) is obtained by specializing Pitman and Yor [43, Proposition 20] to $\alpha = \frac{1}{2}$ and $\theta = 0$. See also Knight [30] for related results. \square

The following result can be read from Maisonneuve [34, Section 8] and Bolthausen [6]. An alternative approach was provided by Greenwood and Pitman [21], and Pitman [39, Sections 4 and 5]. See also Biane and Yor [5, Theorem 6.1], and Revuz and Yor [45, Exercise 4.18, Chapter XII].

Theorem 2.4. [34, 6] *The process*

$$(B - \underline{B})[J, J + 1] := (B_{J+t} - \underline{B}_{J+t}; 0 \leq t \leq 1)$$

is a Brownian meander of length 1, independent of $(B_u; 0 \leq u \leq J)$.

3. THE ARGMIN PROCESS OF BROWNIAN MOTION

In this section, we study the argmin process α of Brownian motion defined by (1.1). In Section 3.1, we deal with the sample path properties of α . In Section 3.2, we provide a conceptual proof that the argmin process α is a Markov process with the Feller property. In Section 3.3, we study the jumps of α by means of a Lévy system. In Section 3.4, we compute the transition kernel of α , and prove Theorem 1.2. In Section 3.5, we derive the infinitesimal generator of α . Finally in Section 3.6, we explain why Dynkin's criterion and the Rogers-Pitman criterion do not apply to the argmin process α .

3.1. Sample path properties. We have mentioned in the introduction that the argmin process $(\alpha_t; t \geq 0)$ takes values in $[0, 1]$, and drifts down at unit speed except for positive jumps. More precisely, we provide the following proposition.

Proposition 3.1. *Let $(\alpha_t; t \geq 0)$ be the argmin process of Brownian motion. Then a.s.*

- (1) $\alpha_t \in [0, 1]$ for all $t \geq 0$, and $(t + \alpha_t; t \geq 0)$ is increasing;
- (2) $(\alpha_t; t \geq 0)$ decreases at unit speed except for
 - (i). jumps from 0 to some $x \in (0, 1)$;
 - (ii). jumps from some $x \in (0, 1)$ to 1.

Proof. (1) The fact $\alpha_t \in [0, 1]$ is straightforward from the definition. Let $0 \leq t < t'$.

- If $t' > t + \alpha_t$, then $t' + \alpha_{t'} > t + \alpha_t$.
- If $t' \leq t + \alpha_t$, then $B_{t+\alpha_t} \leq B_u$ for all $u \in [t', t + \alpha_t]$. This implies that $\alpha_{t'} \geq t + \alpha_t - t'$.



FIGURE 4. LEFT: A jump from 0 to some $x \in (0, 1)$ in the argmin process. RIGHT: A jump from some $x \in (0, 1)$ to 1 in the argmin process.

(2) Observe that $(\alpha_t; t \geq 0)$ is a càdlàg process with only positive jumps. We first check (i). If $\alpha_{t-} = 0$ for some $t > 0$, then $B_u \geq B_t$ for all $u \in [t, t+1]$. We distinguish two cases. In the first case, $B_u > B_t$ for all $u \in [t, t+1]$, which implies that $\alpha_t = 0$. In the second case, $B_u = B_t$ for some $u \in (t, t+1]$, and let $x := \sup\{u \in (0, 1] : B_{t+u} = B_t\}$. If $x = 1$, then there exists an excursion of length 1 in Brownian motion by a space-time shift. But this is excluded by Pitman and Tang [42, Theorem 4]. Thus, $\alpha_t = x \in (0, 1)$.

It remains to check (ii). If $\alpha_{t-} = x \in (0, 1)$ for some $t > 0$, then $B_{t+x} < B_u$ for all $u \in (t+x, t+1)$. We also distinguish two cases. In the first case, $B_{t+1} > B_{t+x}$ which implies that $\alpha_t = x$. In the second case, $B_{t+1} = B_{t+x}$ which yields $\alpha_t = 1$. \square

Next we prove a time reversal property of the argmin process α . By convention, $\alpha_{0-} = \alpha_0$.

Proposition 3.2. *For each fixed $T > 0$, $(1 - \alpha_{(T-t)-}; 0 \leq t \leq T)$ has the same distribution as $(\alpha_t; 0 \leq t \leq T)$.*

Proof. Observe that $(1 - \alpha_{(T-t)-}; 0 \leq t \leq T)$ is also a càdlàg process. Let

$$\tilde{B} := (B_{T+1-u} - B_{T+1}; 0 \leq u \leq T+1) \stackrel{(d)}{=} (B_u; 0 \leq u \leq T+1).$$

Let $\tilde{\alpha}$ be the argmin process of \tilde{B} on $[0, T]$. Hence,

$$(1 - \alpha_{(T-t)-}; 0 \leq t \leq T) = (\tilde{\alpha}_t; 0 \leq t \leq T) \stackrel{(d)}{=} (\alpha_t; 0 \leq t \leq T).$$

\square

Remark 3.3. Let $T > 0$. Note that $\alpha_{T-t} = \alpha_{(T-t)-}$ a.s. for each fixed $t \leq T$, but not simultaneously for all $t \leq T$. This implies that the càglàd $(1 - \alpha_{T-t}; 0 \leq t \leq T)$ has the same finite dimensional distributions as the càdlàg $(\alpha_t; 0 \leq t \leq T)$.

3.2. Markov and Feller property. We provide a soft argument to prove that $(\alpha_t; t \geq 0)$ is a Markov process, and enjoys the Feller property.

For each $t \geq 0$, let

$$\mathcal{G}_t := \sigma(B_s; 0 \leq s \leq t + \alpha_t), \tag{3.1}$$

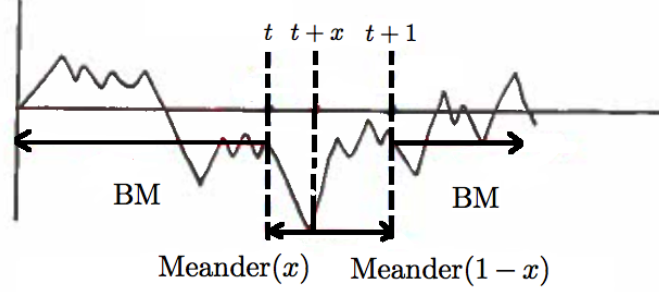
be the σ -field generated by the path B killed at time $t + \alpha_t$. See Chung [8], Pittenger and Shih [44], and Millar [37] for discussion on σ -fields associated to general random times. We do not pursue this direction here.

By Proposition 3.1 (1), $t \mapsto t + \alpha_t$ is increasing. It is not hard to see that for any $s < t$, $s + \alpha_s$ is a measurable function of the path B killed at $t + \alpha_t$. So $(\mathcal{G}_t)_{t \geq 0}$ is a filtration. Now we show that

Proposition 3.4. *The argmin process $(\alpha_t; t \geq 0)$ is time-homogeneous Markov with respect to $(\mathcal{G}_t)_{t \geq 0}$.*

Proof. Fix $t > 0$. By Denisov's decomposition (Theorem 2.1), given $\alpha_t = x$ (i.e. the minimum of B on $[t, t+1]$ is attained at $t+x$), the Brownian path is decomposed into four independent components:

- $(B_{t-s} - B_t; 0 \leq s \leq t)$ is Brownian motion of length t ;
- $(B_{t+x-s} - B_{t+x}; 0 \leq s \leq x)$ is a Brownian meander of length x ;
- $(B_{t+x+s} - B_{t+x}; 0 \leq s \leq 1-x)$ is a Brownian meander of length $1-x$;
- $(B_{t+1+s} - B_{t+1}; s \geq 0)$ is Brownian motion running forever.


 FIGURE 5. Decomposition of Brownian motion given $\alpha_t = x$.

As a consequence, $(B_{t+x+s} - B_{t+x}; s \geq 0)$ and $(B_{t+x-s} - B_{t+x}; 0 \leq s \leq t+x)$ are conditionally independent. Observe that given $\alpha_t = x$,

- for $s > t$, the minimum of B on $[s, s+1]$ cannot be attained on $[s, t+x)$. So $(\alpha_s; s > t)$ is entirely determined by the path $(B_{t+x+s} - B_{t+x}; s \geq 0)$.
- for $s < t$, the minimum of B on $[s, s+1]$ cannot be attained on $(t+x, s+1]$. So $(\alpha_s; s < t)$ is entirely determined by the path $(B_{t+x-s} - B_{t+x}; 0 \leq s \leq t+x)$.

These observations imply that $(\alpha_t; t \geq 0)$ is Markov relative to $(\mathcal{G}_t)_{t \geq 0}$. The time-homogeneity follows from the fact that given $\alpha_t = x$, the law of $(B_{t+x+s} - B_{t+x}; s \geq 0)$ does not involve the time parameter t . \square

We now investigate the Feller property of the argmin process $(\alpha_t; t \geq 0)$. Recall the definition of \mathbf{P}^x from (2.3). Let

$$\alpha_t^x := \alpha_t(B) \quad \text{for } t \geq 0, \quad (3.2)$$

be the argmin process of $(B_t; t \geq 0)$ under \mathbf{P}^x , which makes $\alpha_0 = x \in [0, 1]$. By Denisov's decomposition (Theorem 2.1), for all $f : \mathcal{C}[0, \infty) \rightarrow \mathbb{R}$ bounded and continuous,

$$\mathbb{E}^{\mathbf{W}} f(\alpha_t; t \geq 0) = \int_0^1 \frac{dx}{\pi \sqrt{x(1-x)}} \mathbb{E} f(\alpha_t^x; t \geq 0).$$

where $\mathbb{E}^{\mathbf{W}}$ is the expectation relative to \mathbf{W} .

The Feller property of $(\alpha_t; t \geq 0)$ follows from a direct computation of the transition semigroup $Q_t(x, \cdot)$ of $(\alpha_t^x; t \geq 0)$, which will be given in Section 3.4. But here we provide a conceptual proof.

Proposition 3.5. *The argmin process $(\alpha_t; t \geq 0)$ enjoys the Feller property, and is a strong Markov process.*

Proof. According to Kallenberg [27, Lemma 19.3], it suffices to show that

- (1) for each $t \geq 0$, $\alpha_t^x \rightarrow \alpha_t^y$ in distribution as $x \rightarrow y$;
- (2) for each $x \in [0, 1]$, $\alpha_t^x \rightarrow x$ in probability as $t \rightarrow 0$.

We first prove (1). For $t \geq 1$, α_t^x and α_t^y are both arcsine distributed regardless of $x, y \in [0, 1]$. Consider the case $t \in (0, 1)$. By Denisov's decomposition (Theorem 2.1), for all $f : \mathcal{C}[0, 2] \rightarrow \mathbb{R}$

bounded and continuous,

$$\mathbb{E}^{\mathbf{W}} f(B_u; 0 \leq u \leq 2) = \int_0^1 \frac{dx}{\pi \sqrt{x(1-x)}} \mathbb{E}^{\mathbf{P}^x} f(B_u; 0 \leq u \leq 2).$$

By the explicit scaling construction of Brownian meanders, the law of $(B_u; 0 \leq u \leq 2)$ under \mathbf{P}^x is weakly continuous in x . Moreover, for each $t \in (0, 1)$, $\alpha_t(w)$ is an a.s. continuous functional of $(w_u; 0 \leq u \leq 2)$, from which follows (1).

It remains to prove (2). Observe that for $t < 1$, $1 - t \leq \alpha_t^1 \leq 1$. So (2) is proved in case of $x = 1$. For $x \in [0, 1)$, by Denisov's decomposition (Theorem 2.1),

$$\inf\{u \geq 1; w_u \leq w_x\} > 1 \quad \mathbf{P}^x\text{-a.s.}$$

Therefore, $\mathbb{P}(\alpha_t^x = x - t \text{ for } t \text{ close to } 0) = 1$, which leads to the desired result. \square

3.3. Jumps and Lévy system. We study the jumps of the argmin process α . Recall the definitions of a $\frac{1}{2}$ -stable subordinator, and the age process of a $\frac{1}{2}$ -stable subordinator from (2.5) and (2.7). We begin with the following observation.

Lemma 3.6. *Let $\rho := \inf\{t > 0; \alpha_t = 1\}$ and $\tau := \inf\{t > \rho; \alpha_t = 0\}$. Then $(1 - \alpha_{\rho+t}; 0 \leq t \leq \tau - \rho)$ has the same distribution as the age process of a $\frac{1}{2}$ -stable subordinator until the age first reaches 1.*

Proof. By the strong Markov property of Brownian motion, $(B_{\rho+1+u} - B_{\rho+1}; u \geq 0)$ is still Brownian motion. It is not hard to see that $(1 - \alpha_{\rho+t}; 0 \leq t \leq \tau - \rho)$ is the age process derived from excursions above the past minimum of $(B_{\rho+1+u} - B_{\rho+1}; u \geq 0)$ until this post- $(\rho + 1)$ Brownian motion escapes its past minimum by time 1. This yields the desired result. \square

By Lemma 3.6, let $(l_t^1; t \geq 0)$ be the local times of α at level 1, normalized to match the $\frac{1}{2}$ -stable subordinator. By time-reversal of α (Proposition 3.2), define similarly $(l_t^0; t \geq 0)$ to be the local times of α at level 0. By stationarity of α ,

$$\mathbb{E} l_t^1 / t = \mathbb{E} l_t^0 / t = c \quad \text{for all } t > 0.$$

We will prove in Corollary 4.9 that the constant $c = 1/\sqrt{2\pi}$. These stationary local times also appeared in the work of Leuridan [32].

Before proceeding further, we need the following terminology. Let $(X_t; t \geq 0)$ be a Hunt process on a suitably nice state space E , e.g. locally compact and separable metric space. The pair (Π, C) constituted of a kernel Π on E and a continuous additive functional C is said to be a *Lévy system* for X if for all bounded and measurable function f on $E \times E$,

$$\mathbb{E} \left(\sum_{0 < s \leq t} f(X_{s-}, X_s) 1_{\{X_{s-} \neq X_s\}} \right) = \mathbb{E} \left(\int_0^t dC_s \int_E \Pi(X_{s-}, dy) f(X_{s-}, y) \right). \quad (3.3)$$

The kernel Π is called the *Lévy measure* of the additive functional C . The notion of a Lévy system was formulated by Watanabe [50], the existence of which was proved for a Hunt process under additional assumptions. The proof was simplified by Beneviste and Jacod [3]. See also Meyer [35], Pitman [39] and Sharpe [47, Chapter VIII] for development.

By Proposition 3.5, the argmin process $(\alpha_t; t \geq 0)$ is a Hunt process. Also define a continuous additive functional C by

$$C_t = t + l_t^0 \quad \text{for } t \geq 0. \quad (3.4)$$

The main result is stated as follows, the proof of which relies on Lemmas 3.8 and 3.10.

Theorem 3.7. *Let $(\alpha_t; t \geq 0)$ be the argmin process of Brownian motion, and $(C_t; t \geq 0)$ be the additive functional as in (3.4). Define a kernel Π on $[0, 1]$ by*

$$\Pi(x, dy) = \begin{cases} \Pi^{0\uparrow}(dy) & \text{for } x = 0, \\ \mu^{\uparrow 1}(x)\delta_1 & \text{for } x \in (0, 1), \\ 0 & \text{for } x = 1. \end{cases} \quad (3.5)$$

where δ_1 is the point mass at 1,

$$\Pi^{0\uparrow}(dy) := \frac{dy}{\sqrt{2\pi y^3(1-y)}} \quad \text{for } 0 < y < 1, \quad (3.6)$$

and

$$\mu^{\uparrow 1}(x) := \frac{1}{2(1-x)} \quad \text{for } 0 < x < 1. \quad (3.7)$$

Then (C, Π) is a Lévy system of $(\alpha_t; t \geq 0)$.

Recall from Proposition 3.1 (2) that $(\alpha_t; t \geq 0)$ can only have (i). jumps from 0 to some $x \in (0, 1)$, and (ii). jumps from some $x \in (0, 1)$ to 1. We start by computing the jump rate of α from $x \in (0, 1)$ to 1.

Lemma 3.8. *Let $(\alpha_t; t \geq 0)$ be the argmin process of Brownian motion.*

(1) *Let $x > y \geq 0$. The probability that α^x decreases at unit speed from x to y with no jumps is given by*

$$s(x, y) = \sqrt{\frac{1-x}{1-y}}. \quad (3.8)$$

(2) *For $x \in (0, 1)$, the jump rate of α per unit time from x to 1 is $\mu^{\uparrow 1}(x)$ defined by (3.7).*

Proof. (1) Let $\tilde{x} := 1 - x$ and $\tilde{y} := 1 - y$. By Denisov's decomposition (Theorem 2.1), under \mathbf{P}^x , $(B_{x+t} - B_x; 0 \leq t \leq \tilde{x})$ is a Brownian meander of length \tilde{x} , independent of Brownian motion $(B_{1+t} - B_1; t \geq 0)$.

Let Z_{x-y} be normally distributed with mean 0 and variance $x - y$. Let $R^{\tilde{x}}$ be Rayleigh distributed with parameter \tilde{x} , independent of Z_{x-y} . We have

$$\begin{aligned} s(x, y) &= \mathbf{P}^x \left(\inf_{t \leq x-y} (B_{1+t} - B_1) > B_x - B_1 \right) \\ &= \mathbb{P}(|Z_{x-y}| < R^{\tilde{x}}) \\ &= 2 \int_0^\infty \mathbb{P}(R^{\tilde{x}} > z) \cdot \frac{1}{\sqrt{2\pi(x-y)}} \exp\left(-\frac{z^2}{2(x-y)}\right) dz = \sqrt{\frac{\tilde{x}}{\tilde{y}}}, \end{aligned}$$

where the second equality follows from the reflection principle of Brownian motion, and the fact that a Brownian meander of length \tilde{x} evaluated at time \tilde{x} is Rayleigh distributed with parameter \tilde{x} , whose density is given by (2.1).

(2) Note that

$$s(x, y) := \exp\left(-\int_y^x dz \mu^{\uparrow 1}(z)\right).$$

We obtain the the jump rate (3.7) by taking derivative of (3.8) with respect to x . \square

Remark 3.9. We provide an alternative approach to Lemma 3.8. Consider the excursions above the past minimum of $(B_t; t \geq 0)$. Given $\alpha_0 = x$, it must be a ladder time; that is the starting time of an excursion. Thus, the probability that α jumps to 1 on $(0, dt]$ given $\alpha_0 = x$ is the same as that of an excursion terminates in dt given that it has reached length \tilde{x} .

Let ζ be the length of such an excursion. By (2.6), the aforementioned probability is given by

$$\mathbf{n}(\zeta \in \tilde{x} + dt | \zeta > \tilde{x}) = \frac{\Lambda(d\tilde{x})/d\tilde{x}}{\Lambda(\tilde{x}, \infty)} dt = \frac{1}{2\tilde{x}} dt,$$

where $\Lambda(dx)$ is the Lévy measure of a $\frac{1}{2}$ -stable subordinator as in (2.5). This gives the jump rate (3.7).

To conclude this subsection, we compute the Lévy measure of jumps of α in from 0.

Lemma 3.10. *Let $(\alpha_t; t \geq 0)$ be the argmin process of Brownian motion. For $y \in (0, 1)$, the Lévy measure of jumps of α per unit local time in from 0 is $\Pi^{0\uparrow}(dy)$ defined by (3.6).*

Proof. On one hand, the mean number of jumps per unit time from 0 to dy near y is $\Pi^{0\uparrow}(dy)\mathbb{E}l_1^0 = \Pi^{0\uparrow}(dy)/\sqrt{2\pi}$. On the other hand, the mean number of jumps per unit time from dy near $1 - y$ to 1 is given by

$$\frac{dy}{\pi\sqrt{y(1-y)}}\mu^{\uparrow 1}(1-y) = \frac{dy}{2\pi\sqrt{y^3(1-y)}}.$$

By Proposition 3.2, we identify these two quantities and obtain the Lévy measure (3.6). \square

3.4. Transition kernel. We complete the proof of Theorem 1.2. Recall the definition of $(\alpha_t^x; t \geq 0)$ from (3.2), which is viewed as the argmin process α conditioned on $\alpha_0 = x$. For $0 \leq b \leq 1$, let

$$\tau_b^x := \inf\{t > 0; \alpha_t^x = b\}$$

to be the first time at which $(\alpha_t^x; t \geq 0)$ hits level b . Also recall the definition of $\mu^{\uparrow 1}(x)$ from (3.7). We start with a lemma whose proof is straightforward.

Lemma 3.11. *For $0 < x < 1$,*

$$\mathbb{P}(\tau_1^x \in dt) = \mu^{\uparrow 1}(x-t)s(x, x-t)dt \quad \text{if } 0 < t < x, \quad (3.9)$$

where $\mu^{\uparrow 1}(x)$ is given by (3.7) and $s(x, y)$ is given by (3.8).

Proof of Theorem 1.2. The first part of Theorem 1.2 has been proved as Proposition 3.5. Now we compute the transition kernel $Q_t(x, dy)$ for $t > 0$ and $x \in [0, 1]$ of $(\alpha_t; t \geq 0)$.

Observe that α_{t+s} and α_s are independent for all $t \geq 1$. By Proposition 1.1,

$$Q_t(x, dy) = \frac{1_{\{0 < y < 1\}}}{\pi\sqrt{y(1-y)}} dy \quad \text{for } t \geq 1 \text{ and } x \in [0, 1], \quad (3.10)$$

which is the invariant measure of the argmin process α .

Given $\alpha_0 = 1$, we have $B_u \geq B_1$ for all $u \in [0, 1]$. So for $0 < t \leq 1$,

$$t + \alpha_t^1 \stackrel{(d)}{=} \sup \left\{ s \in [1, 1+t]; B_s = \max_{u \in [1, 1+t]} B_u \right\}.$$

Consequently, α_t^1 is the arcsine distribution rescaled linearly into $[1-t, 1]$. That is,

$$Q_t(1, dy) = \frac{1_{\{1-t < y < 1\}}}{\pi \sqrt{(1-y)(y+t-1)}} dy \quad \text{for } 0 < t \leq 1. \quad (3.11)$$

By conditioning on τ_1^x with $\tau_1^x \leq x$, we have for $0 < t \leq x \leq 1$,

$$Q_t(x, dy) = s(x, x-t) \delta_{x-t}(dy) + \int_{x-t}^x dz \mu^{\uparrow 1}(z) s(x, z) Q_{t+z-x}(1, dy), \quad (3.12)$$

while for $0 < x < t \leq 1$,

$$Q_t(x, dy) = s(x, 0) P_{t-x}(0, dy) + \int_0^x dz \mu^{\uparrow 1}(z) s(x, z) Q_{t+z-x}(1, dy). \quad (3.13)$$

In the case $t \leq x$ there is an atom of probability $s(x, x-t)$ at $x-t$, whereas in the case $t > x$ this atom is replaced by probability $s(x, 0)$ redistributed according to $Q_{t-x}(0, dy)$. For $t = 1$, we know that $Q_1(x, dy)$ is arcsine distributed, whatever x . So this case gives a formula for $Q_u(0, dy)$ for any $0 < u < 1$ with $u := 1 - x$. That is,

$$Q_u(0, dy) = \frac{1}{s(1-u, 0)} \left[\frac{1_{\{0 < y < 1\}} dy}{\pi \sqrt{y(1-y)}} - \int_0^{1-u} dz \mu^{\uparrow 1}(z) s(1-u, z) Q_{u+z}(1, dy) \right]. \quad (3.14)$$

It remains to evaluate the r.h.s. of (3.12)-(3.14). By (3.8) and (3.11), we get

$$s(x, x-t) = \sqrt{\frac{1-x}{1-x+t}},$$

and

$$\begin{aligned} & \int_{x-t}^x dz \mu^{\uparrow 1}(z) s(x, z) Q_{t+z-x}(1, dy) \\ &= \int_{x-t}^x \frac{dz}{2(1-z)} \cdot \sqrt{\frac{1-x}{1-z}} \cdot \frac{1_{\{y \geq 1+x-t-z\}} dy}{\pi \sqrt{(1-y)(y+t+z-x-1)}} \\ &= \frac{dy \sqrt{1-x}}{2\pi \sqrt{1-y}} \int_{1-t+x-y}^x \frac{1_{\{y \geq 1-t\}} dz}{\sqrt{(1-z)^3(y+t+z-x-1)}} \\ &= \frac{dy \sqrt{1-x}}{\pi(y+t-x) \sqrt{1-y}} \left[\sqrt{\frac{(y+t+z-x-1)^+}{1-z}} \right]_{1-t+x-y}^x \\ &= \frac{\sqrt{(y+t-1)^+}}{\pi(y+t-x) \sqrt{1-y}} dy. \end{aligned}$$

By injecting these expressions into (3.12), we obtain

$$Q_t(x, dy) = \sqrt{\frac{1-x}{1-x+t}} \delta_{x-t}(dy) + \frac{\sqrt{(y+t-1)^+}}{\pi(y+t-x) \sqrt{1-y}} dy \quad \text{for } t \leq x \leq 1. \quad (3.15)$$

Similarly, we get from (3.13) that for $x < t \leq 1$,

$$Q_t(x, dy) = \sqrt{1-x} Q_{t-x}(0, dy) + \frac{dy \sqrt{1-x}}{\pi(y+t-x) \sqrt{1-y}} \left[\sqrt{\frac{(y+t-1)^+}{1-x}} - \sqrt{(y+t-1-x)^+} \right];$$

and from (3.14) that for $0 < u < 1$,

$$Q_u(0, dy) = \frac{\sqrt{u} + \sqrt{y(y+u-1)^+}}{\pi(y+u)\sqrt{y(1-y)}} dy.$$

By combining the above expressions, we obtain

$$Q_t(x, dy) = \frac{\sqrt{(1-x)(t-x)} + \sqrt{y(y+t-1)^+}}{\pi(y+t-x)\sqrt{y(1-y)}} dy \quad \text{for } x < t \leq 1. \quad (3.16)$$

□

3.5. Infinitesimal generator. We consider the infinitesimal generator L of the argmin process $(\alpha_t; t \geq 0)$. For each $f \in \mathcal{C}[0, 1]$ and $x \in [0, 1]$, let

$$Lf(x) := \lim_{t \downarrow 0} \frac{Q_t f(x) - f(x)}{t}, \quad \text{if the limit exists.}$$

The domain of the generator $\mathcal{D}(L)$ is the set of $f \in \mathcal{C}[0, 1]$ for which the limit exists for all $x \in [0, 1]$. Though we are not able to identify the domain of L , we exhibit a large subset S of $\mathcal{D}(L)$, given by (3.26).

The following theorem gives the infinitesimal generator of the argmin process $(\alpha_t; t \geq 0)$.

Theorem 3.12. *Let L be the infinitesimal generator of $(\alpha_t; t \geq 0)$, and $\mathcal{D}(L)$ be the domain of L . If $f \in \mathcal{D}(L) \cap \mathcal{C}^1[0, 1]$, then*

$$Lf(x) = \begin{cases} -f'(x) + \frac{f(1) - f(x)}{2(1-x)} & \text{for } 0 \leq x < 1, \\ -\frac{1}{2}f'(1) & \text{for } x = 1. \end{cases} \quad (3.17)$$

Let \mathbf{P}^x be defined by (2.3). Then \mathbf{P}^x is the unique probability measure such that

$$\alpha_0 = x \quad \mathbf{P}^x\text{-a.s.} \quad (3.18)$$

and for each $f \in \mathcal{D}(L) \cap \mathcal{C}^1[0, 1]$,

$$\left(M_t^f := f(\alpha_t) - f(\alpha_0) - \int_0^t Lf(\alpha_s) ds \right)_{t \geq 0} \quad \text{is a martingale under } \mathbf{P}^x. \quad (3.19)$$

Proof. For $0 < t \leq x < 1$, we obtain by (1.4) that

$$\begin{aligned} Q_t f(x) &= \sqrt{\frac{1-x}{1-x+t}} f(x-t) + \int_{1-t}^1 \frac{\sqrt{y+t-1}}{\pi\sqrt{1-y}(y+t-x)} f(y) dy \\ &= \sqrt{\frac{1-x}{1-x+t}} f(x-t) + t \int_0^1 \frac{\sqrt{y}}{\pi\sqrt{1-y}(ty+1-x)} f(ty+1-t) dy \\ &= f(x) - t \left(f'(x) + \frac{f(x)}{2(1-x)} \right) + \frac{tf(1)}{1-x} \int_0^1 \frac{\sqrt{y} dy}{\pi\sqrt{1-y}} + o(t) \\ &= f(x) - t \left(f'(x) + \frac{f(x) - f(1)}{2(1-x)} \right) + o(t). \end{aligned} \quad (3.20)$$

Consequently,

$$Lf(x) = -f'(x) + \frac{f(1) - f(x)}{2(1-x)} \quad \text{for } 0 < x < 1. \quad (3.21)$$

Note that $x \mapsto Lf(x)$ is continuous on $[0, 1]$. By taking $x \rightarrow 0$ and $x \rightarrow 1$ in (3.21), we get

$$Lf(0) = -f'(0) + \frac{f(1) - f(0)}{2} \quad \text{and} \quad Lf(1) = -\frac{1}{2}f'(1). \quad (3.22)$$

The second part follows from general theory of the martingale problem for Markov processes, see Either and Kurtz [17, Chapter 4]. \square

For $0 = x < t \leq 1$, we obtain by (1.4) that

$$\begin{aligned} Q_t f(0) &= \sqrt{t} \int_0^1 \frac{f(y)}{\pi(y+t)\sqrt{y(1-y)}} dy + \int_{1-t}^1 \frac{\sqrt{y+t-1}}{\pi(y+t)\sqrt{1-y}} f(y) dy \\ &= \sqrt{t} \int_0^1 \frac{f(y)}{\pi(y+t)\sqrt{y(1-y)}} dy + t \frac{f(1)}{2} + o(t). \end{aligned} \quad (3.23)$$

Now assume $f \in \mathcal{C}^1[0, 1]$ so that

$$\begin{aligned} \int_0^1 \frac{f(y)}{\pi(y+t)\sqrt{y(1-y)}} dy &= f(0) \int_0^1 \frac{1}{\pi(y+t)\sqrt{y(1-y)}} dy \\ &\quad + f'(0) \int_0^1 \frac{y}{\pi(y+t)\sqrt{y(1-y)}} dy + \int_0^1 \frac{f(y) - f(0) - yf'(0)}{\pi(y+t)\sqrt{y(1-y)}} dy \end{aligned}$$

Observe that

$$\begin{aligned} \int_0^1 \frac{1}{\pi(y+t)\sqrt{y(1-y)}} dy &= \frac{1}{\sqrt{t(t+1)}} = \frac{1}{\sqrt{t}} - \frac{1}{2}\sqrt{t} + \mathcal{O}(t^{\frac{3}{2}}), \\ \int_0^1 \frac{y}{\pi(y+t)\sqrt{y(1-y)}} dy &= 1 - \sqrt{t} + \mathcal{O}(t^{\frac{3}{2}}), \end{aligned}$$

and

$$\int_0^1 \frac{f(y) - f(0) - yf'(0)}{\pi(y+t)\sqrt{y(1-y)}} dy = \int_0^1 \frac{f(y) - f(0)}{\pi\sqrt{y^3(1-y)}} dy - f'(0) + \mathcal{O}(t^{\frac{3}{2}}).$$

By injecting these expressions into (3.23), we get

$$Q_t f(0) = f(0) + \sqrt{t} \int_0^1 \frac{f(y) - f(0)}{\pi\sqrt{y^3(1-y)}} dy + t \left(-f'(0) + \frac{f(1) - f(0)}{2} \right) + o(t). \quad (3.24)$$

Note that

$$\frac{Q_t f(0) - f(0)}{t} \rightarrow -f'(0) + \frac{f(1) - f(0)}{2} \quad \text{as } t \rightarrow 0.$$

So the coefficient of \sqrt{t} in (3.24) is forced to be 0. That is,

$$\int_0^1 \frac{f(y) - f(0)}{\pi\sqrt{y^3(1-y)}} dy = 0 \quad (3.25)$$

Therefore,

$$\mathcal{S} := \left\{ f \in \mathcal{C}^1[0, 1] : \int_0^1 \frac{f(y) - f(0)}{\pi\sqrt{y^3(1-y)}} dy = 0 \right\} \subset \mathcal{D}(L). \quad (3.26)$$

If $f(x) = \sum_{n=0}^{\infty} a_n x^n$ for $0 \leq x \leq 1$, then f belongs to \mathcal{S} if and only if its coefficients $(a_n; n \geq 0)$ satisfy

$$\sum_{n=1}^{\infty} \binom{2n-2}{n-1} \frac{a_n}{2^{2n-2}} = 0.$$

But we do not know whether the domain $\mathcal{D}(L)$ is exactly \mathcal{S} or not. We leave this open for further investigation.

3.6. Breakdown of Dynkin's and Rogers-Pitman criterion. In this part, we explain why Dynkin's criterion, and the Rogers-Pitman criterion fail to prove that $(\alpha_t; t \geq 0)$ is Markov. Before proceeding further, we recall these sufficient conditions for a function of a Markov process to be Markov.

Let $(X_t; t \geq 0)$ be a continuous-time Markov process on a measurable state space (E, \mathcal{E}) , with initial distribution λ and transition semigroup $(P_t; t \geq 0)$. Let (E', \mathcal{E}') be a second measurable space, and $\phi : (E, \mathcal{E}) \rightarrow (E', \mathcal{E}')$ be a measurable function.

Dynkin [16] initiated the study of Markov functions, and gave a condition for $(\phi(X_t); t \geq 0)$ to be Markov for all initial distributions λ . Later Rogers and Pitman [46] made a simple observation: if there exists a Markov kernel $\Lambda : E' \times \mathcal{E} \ni (y, A) \mapsto \Lambda(y, A) \in \mathbb{R}_+$ such that for all $t \geq 0$ and $A \in \mathcal{E}$,

$$\mathbb{P}(X_t \in A | \phi(X_s), 0 \leq s \leq t) = \Lambda(\phi(X_t), A) \quad a.s., \quad (3.27)$$

then $(\phi(X_t); t \geq 0)$ is Markov with transition kernels

$$Q_t = \Lambda P_t \Phi \quad \text{for all } t \geq 0,$$

where Φ is the Markov kernel from E' to E induced by ϕ : $\Phi(x, B) = \delta_{\phi(x)}(B)$ for $x \in E$ and $B \in \mathcal{E}'$. The following theorem provides a sufficient condition for (3.27) to hold.

Theorem 3.13 (Rogers-Pitman criterion). [46] *Let Φ be derived from $\phi : E \rightarrow E'$ as above. Assume that there exists a Markov kernel Λ from E' to E such that*

- (i). $\Lambda \Phi = I$, the identity kernel;
- (ii). for each $t \geq 0$, the Markov kernel $Q_t = \Lambda P_t \Phi$ satisfies the intertwining relation $\Lambda P_t = Q_t \Lambda$.

Let $(X_t); t \geq 0$ be Markov with initial distribution $\lambda = \Lambda(y, \cdot)$ for some $y \in E'$ and semigroup $(P_t; t \geq 0)$. Then (3.27) holds, and $(\phi(X_t); t \geq 0)$ is Markov with starting state y and transition semigroup $(Q_t; t \geq 0)$.

Note that if instead of (ii),

$$P_t \Phi = \Phi Q_t \quad \text{for all } t \geq 0,$$

for a Markov kernel Q_t on E' , then $(\phi(X_t); t \geq 0)$ is Markov for all initial distributions λ . This recovers Dynkin's criterion [16].

As shown by (1.2), the argmin process α is a measurable function of the space-time shift process $(\Theta_t; t \geq 0)$ whose transition kernel is given by

$$P_t(w, d\tilde{w}) = \delta_{\Theta_t w}(d\tilde{w}), \quad (3.28)$$

where $\Theta_t w := (w_{t+u} - w_t; t \geq 0)$ is the space-time shift of w on $\mathcal{C}[0, \infty)$. For $w \in \mathcal{C}[0, \infty)$ and $t \geq 0$, let

$$\alpha_t(w) := \sup \left\{ s \in [0, 1] : w_{t+s} = \min_{u \in [0, 1]} w_{t+u} \right\}. \quad (3.29)$$

The Markov kernel Φ induced by the argmin function is given by

$$\Phi(w, \cdot) = \delta_{\alpha_0(w)}(\cdot) \quad \text{for all } w \in \mathcal{C}[0, \infty). \quad (3.30)$$

We first show that Dynkin's criterion does not hold.

Proposition 3.14. *Let $(P_t; t \geq 0)$ be the semigroup of the space-time shift process Θ given by (3.28), $(Q_t; t \geq 0)$ be the semigroup of the argmin process α given by (1.4), and Φ be the Markov kernel defined by (3.30). Then*

$$P_t \Phi(w, \cdot) \neq \Phi Q_t(w, \cdot) \quad \text{for all } t > 0.$$

Proof. Observe that for all $t \geq 0$,

$$P_t \Phi(w, \cdot) = \delta_{\alpha_t(w)}(\cdot) \quad \text{and} \quad \Phi Q_t(w, \cdot) = Q_t(\alpha_0(w), \cdot),$$

where $\alpha_t(w)$ is given by (3.29). From this follows the result. \square

Recall the definition of \mathbf{P}^x from (2.3). By Denisov's decomposition (Theorem 2.1), the condition (3.27) amounts to

$$\Lambda(x, \cdot) = \mathbf{P}^x \quad \text{for all } x \in [0, 1]. \quad (3.31)$$

The following result shows that Rogers-Pitman intertwining criterion does not hold.

Proposition 3.15. *Let $(P_t; t \geq 0)$ be the semigroup of the space-time shift process Θ given by (3.28), $(Q_t; t \geq 0)$ be the semigroup of the argmin process α given by (1.4), and Λ be the Markov kernel defined by (3.31). Then for each $t \in (0, 1]$,*

$$\Lambda P_t(t, \cdot) \neq Q_t \Lambda(t, \cdot).$$

Proof. Observe that $\Lambda P_t(t, \cdot) = \overrightarrow{\mathbf{M}^{1-t}} \otimes \mathbf{W}$: the law of a Brownian meander of length $1 - t$ concatenated by Brownian motion. Thus,

$$\mathbb{E}^{\Lambda P_t(t, \cdot)}[w_1] = \sqrt{\frac{\pi}{2}(1-t)}.$$

Next by (1.4),

$$Q_t \Lambda(t, \cdot) = \int Q_t(t, dy) \mathbf{P}^y = \sqrt{1-t} \mathbf{P}^0 + \int_{1-t}^1 \frac{\sqrt{y+t-1}}{\pi y \sqrt{1-y}} \mathbf{P}^y dy,$$

which implies that

$$\begin{aligned} \mathbb{E}^{Q_t \Lambda(t, \cdot)}[w_1] &= \sqrt{\frac{\pi}{2}(1-t)} + \frac{1}{\sqrt{2\pi}} \int_{1-t}^1 \frac{\sqrt{y+t-1}}{\sqrt{y(1-y)}} dy \\ &> \mathbb{E}^{\Lambda P_t(t, \cdot)}[w_1] \quad \text{for } t \in (0, 1]. \end{aligned}$$

This yields the desired result. \square

Remark 3.16. Let $(X_t; t \geq 0)$ be the $\mathcal{C}[0, 1]$ -valued moving-window process of Brownian motion B defined by

$$X_t := (B_{t+u} - B_t; 0 \leq u \leq 1) \quad \text{for all } t \geq 0.$$

So X is the projection of Θ from $\mathcal{C}[0, \infty)$ to $\mathcal{C}[0, 1]$. The fact that X is a Markov process follows from applying the Rogers-Pitman criterion with

$$\Lambda(w, d\tilde{w}) = 1(\tilde{w} = w \otimes w') \mathbf{W}(dw'),$$

where \otimes is the usual path concatenation. Note that the argmin process

$$\alpha_t = \hat{\alpha}(X_t) \quad \text{for all } t > 0,$$

where

$$\hat{\alpha}(w) := \sup \left\{ s \in [0, 1] : w_s = \min_{u \in [0, 1]} w_u \right\} \quad \text{for } w \in \mathcal{C}[0, 1].$$

The same argument as in Proposition 3.15 shows that as a function of the $\mathcal{C}[0, 1]$ -valued process X , the argmin process α does not satisfy the Rogers-Pitman criterion.

4. THE (a, b) -MINIMA SET OF BROWNIAN MOTION

In this section, we study the (a, b) -minima set $\mathcal{M}_{a,b}$ of Brownian motion defined by (1.5). In Section 4.1, we consider the renewal property of the set $\mathcal{M}_{a,b}$, and provide an alternative proof of Theorem 1.4. In Section 4.2, we give an explicit construction for times of the set $\mathcal{M}_{1,1}$, which implies Theorem 1.5. Finally in Section 4.3, we deal with the sample path of Brownian motion between two $(1, 1)$ -minima. There Theorem 1.6 is proved.

4.1. Renewal structure of (a, b) -minima. We provide an alternative proof of Theorem 1.4. Recall that the argmin process α is a stationary Markov process, whose

- invariant measure has density $f(x)$, $0 < x < 1$ given by (1.3);
- transition kernel $Q_t(x, \cdot)$, $t > 0$ and $x \in [0, 1]$ is given by (1.4).

For $y \neq x - t$, write $Q_t(x, dy) = q_t(x, y)dy$. The following lemma, which is crucial in Leuridan's proof of Theorem 1.4, can be derived from Proposition 1.1 and Theorem 1.2.

Lemma 4.1. [38, 32] *Given a measurable set $A \subset \mathbb{R}^+$, let $N_{a,b}(A) := \#(\mathcal{M}_{a,b} \cap A)$ be the counting measure of (a, b) -minima. Then*

$$\mathbb{E}(N_{a,b}(dt)) = \frac{dt}{\pi\sqrt{ab}} \tag{4.1}$$

and for $0 \leq s < t$,

$$\mathbb{E}(N_{a,b}(ds)N_{a,b}(dt)) = \frac{1}{\pi\sqrt{ab}}h_{a,b}(t-s)dsdt, \tag{4.2}$$

where $h_{a,b}$ is defined by (1.10).

Proof. Observe that

$$N_{a,b}(dt) = 1 \iff \text{the minimum of } B \text{ on } [t-a, t+b] \text{ is achieved in } dt.$$

By Brownian scaling, the latter has the same probability as that of $\left\{(a+b)\alpha_{\frac{t-a}{a+b}} \in da\right\}$. So

$$\mathbb{E}(N_{a,b}(dt)) = \frac{1}{a+b} f\left(\frac{a}{a+b}\right) dt = \frac{dt}{\pi\sqrt{ab}}.$$

A similar argument shows that

$$\begin{aligned} \mathbb{E}(N_{a,b}(ds)N_{a,b}(dt)) &= \mathbb{E}(N_{a,b}(ds)) \cdot \frac{1}{a+b} q_{\frac{t-s}{a+b}}\left(\frac{a}{a+b}, \frac{a}{a+b}\right) dt \\ &= \frac{1}{\pi\sqrt{ab}} h_{a,b}(t-s) ds dt. \end{aligned}$$

□

Proof of Theorem 1.4. Note that $\mathcal{M}_{a,b} - a := (T_i^{a,b} - a; i \geq 1)$ is a renewal process with stationary delay. By Lemma 4.1, $h_{a,b}(\cdot)$ is the *renewal function* of the point process $\mathcal{M}_{a,b} - a$. The formula (1.11) follows from Daley and Vere-Jones [12, Example 5.4(b)]. By renewal theory, the law of $T_1^{a,b} - a$ is obtained first by size-biasing the inter-arrival time distribution (1.11), and then by stick-breaking uniformly at random, see Thorisson [48]. This gives the formula (1.12). □

4.2. Construction of $(1,1)$ -minima. We consider the case $a = b = 1$ by studying the law of Brownian fragments between $(1,1)$ -minima. Let T_1, T_2, \dots with $0 < T_1 < T_2 < \dots$ be times of $(1,1)$ -minima set of Brownian motion. Now we give a path construction for T_1, T_2, \dots , from which the renewal property of $\mathcal{M}_{1,1}$ is clear. In particular, Theorem 1.5 is a corollary of this construction.

Construction of T_1 Let J be the first descending ladder time of B , from which starts an excursion above the minimum of length exceeding 1. The Laplace transform of J is given by (1.17). By Theorem 2.4, $(B - \underline{B})[J, J+1]$ is a Brownian meander of length 1.

If $J \geq 1$, then $T_1 = J$. If not, we start afresh Brownian motion at the stopping time $J+1$; that is $B^1 := (B_{J+1+t} - B_{J+1}; t \geq 0)$. Let J_1 be constructed as J for B^1 . Thus, $J_1 \in \mathcal{LE}$, and $(B^1 - \underline{B}^1)[J_1, J_1+1]$ is a Brownian meander of length 1. Now we look backward a unit from J_1 to see whether $J_1 \in \mathcal{RE}$ or not. If $J_1 \in \mathcal{RE}$, then $T_1 = J_1$. If not, we start afresh Brownian motion at J_1+1 and proceed as before until a $(1,1)$ -minima is found.

Construction of T_{i+1} given T_i By induction, $(B_{T_i+t} - B_{T_i}; 0 \leq t \leq 1)$ is a Brownian meander of length 1. Now it suffices to start afresh Brownian motion at T_i+1 , and proceed as in the construction of T_1 .

Evaluation of the geometric rate Recall that Δ is distributed as $T_{i+1} - T_i$. Let

$$N := \inf\{i \geq 1 : J_i \in \mathcal{RE}\}. \quad (4.3)$$

It is easy to see that N is geometrically distributed on $\{1, 2, \dots\}$ with parameter $\mathbb{P}(J_1 \in \mathcal{RE})$. Note that J_i depends on the event $\{N = i\}$, but is independent of the event $\{N \geq i\}$. In fact, N is a stopping time of a sequence of i.i.d. path fragments, each starting with a meander and continuing with an independent Brownian motion until time J_i . By Wald's identity,

$$\mathbb{E}\Delta = \mathbb{E}N \cdot (1 + \mathbb{E}J_1) = \frac{1 + \mathbb{E}J_1}{\mathbb{P}(J_1 \in \mathcal{RE})}.$$

Now by (1.14), we get

$$\mathbb{P}(J_1 \in \mathcal{RE}) = \frac{2}{\pi}. \quad (4.4)$$

In view of the dependence of J_i and the event $\{N = i\}$, the evaluation of the geometric rate in the distribution of N is quite indirect. Here is a more direct approach.

Consider the construction of J_1 as J for a copy of Brownian motion preceded by an independent meander of length 1. It is straightforward that

$$\mathbb{P}(J_1 \in \mathcal{RE} \text{ and } J_1 \geq 1) = P(J_1 \geq 1) = 1 - \frac{2}{\pi}, \quad (4.5)$$

where the second equality is obtained by integrating (2.12) over $[0, 1]$. The evaluation of $\mathbb{P}(J_1 \in \mathcal{RE} \text{ and } J_1 < 1)$ is more tricky, which relies on the following lemma.

Lemma 4.2. *Let $(L_t^{br}; 0 \leq t \leq 1)$ be the local time process of a Brownian bridge of length 1 at level 0. Then*

$$\mathbb{P}(L_t^{br} > x) = e^{-\frac{x^2}{2}} \mathbb{P}\left(|B_1| > x\sqrt{\frac{1-t}{t}}\right) \quad \text{for } t \in [0, 1], x > 0. \quad (4.6)$$

Proof. It can be read from Pitman [40, (3)] that for $t \in [0, 1]$, $x > 0$ and $y \in \mathbb{R}$,

$$\mathbb{P}(L_t^{br} > x | B_t^{br} \in dy) = \exp\left(-\frac{1}{2} \left[\left(\frac{|y|}{t} + \frac{x}{t}\right)^2 - \frac{y^2}{t}\right]\right). \quad (4.7)$$

By integrating (4.7) with respect to the normal density of B_t^{br} with mean 0 and variance $t(1-t)$, we obtain (4.6). \square

Let $-\xi$ be the level of the minimum of the free Brownian part of the path at time J_1 so that $J_1 = \sigma_\xi$, where σ is $\frac{1}{2}$ -stable subordinator with jumps of size larger than 1 deleted. Recall from Section 2.2 that ξ is exponentially distributed with parameter $\sqrt{2/\pi}$. By letting $T_x := \inf\{t > 0 : B_t = x\}$, we obtain for $0 < t < 1$,

$$\begin{aligned} \mathbb{P}(\xi \in dx, J_1 \in dt) &= \sqrt{2/\pi} dx \mathbb{P}(T_x \in dt) \\ &= \frac{x}{\pi t^{3/2}} e^{-\frac{x^2}{2t}} dt dx. \end{aligned} \quad (4.8)$$

By time-reversing the Biane-Yor construction [5] of Brownian meander minus its future minimum process (see also Bertoin and Pitman [4, Theorem 3.1]), we get

$$\begin{aligned} \mathbb{P}(J_1 \notin \mathcal{RE} \text{ and } J_1 < 1) &= \int_0^\infty \int_0^1 \mathbb{P}(L_{1-t}^{br} > x) \mathbb{P}(J_1 \in dt, \xi \in dx) \\ &= 1 - \frac{2}{\pi}, \end{aligned} \quad (4.9)$$

where the last equality is obtained by plugging in (4.6) and (4.8). Now (4.4) follows readily from (4.5) and (4.9).

Proof of Theorem 1.5. For any random variable X , let $\Phi_X(\lambda)$ be the Laplace transform of X . The identity (1.18) is clear from the preceding construction. It implies that

$$\begin{aligned}\mathbb{E}T_1 &= \mathbb{E}J + \mathbb{P}(J < 1)\mathbb{E}\Delta \\ &= 1 + \frac{2}{\pi} \cdot \pi = 3,\end{aligned}$$

where the second equality follows from (2.10), (2.12), and (1.14). In addition,

$$\Phi_{T_1}(\lambda) = \mathbb{E}(e^{-\lambda J} 1_{\{J \geq 1\}}) + \Phi_{\Delta}(\lambda) \mathbb{E}(e^{-\lambda J} 1_{\{J < 1\}}). \quad (4.10)$$

By integrating with respect to (2.12), we get

$$\mathbb{E}(e^{-\lambda J} 1_{\{J < 1\}}) = \frac{\operatorname{erf}(\sqrt{\lambda})}{\sqrt{\pi\lambda}}. \quad (4.11)$$

By injecting (4.11) into (4.10), we obtain

$$\Phi_{T_1}(\lambda) = \Phi_J(\lambda) - \frac{\operatorname{erf}(\sqrt{\lambda})}{\sqrt{\pi\lambda}}(1 - \Phi_{\Delta}(\lambda)). \quad (4.12)$$

Recall from (1.9) that $T_1 - 1$ is the stationary delay for a renewal process with inter-arrival time distributed according to Δ . By renewal theory,

$$\mathbb{P}(T_1 - 1 \in dt)/dt = \frac{1}{\pi} \mathbb{P}(\Delta > t), \quad (4.13)$$

which implies that

$$\Phi_{T_1}(\lambda) = \frac{e^{-\lambda}}{\pi\lambda}(1 - \Phi_{\Delta}(\lambda)). \quad (4.14)$$

Combining (4.12) and (4.14) yields

$$\Phi_{T_1}(\lambda) = e^{-\lambda}(\Phi_J(\lambda))^2, \quad (4.15)$$

and

$$\Phi_{\Delta}(\lambda) = 1 - \pi\lambda(\Phi_J(\lambda))^2, \quad (4.16)$$

Now the identity (1.19) follows readily from (4.15). By plugging the formula (1.17) for $\Phi_J(\lambda)$ into (4.15) and (4.16), we get (1.20) and (1.21).

Let H be distributed as $B_{T_{i+1}} - B_{T_i}$, $i \geq 1$. Recall from (2.1) that a Brownian meander evaluated at time 1 has Rayleigh distribution with parameter 1. It is clear from the above construction that

$$H \stackrel{(d)}{=} \sum_{i=1}^N (R_i - \xi_i), \quad (4.17)$$

where $(R_i)_{i \geq 1}$ are i.i.d. Rayleigh distributed with parameter 1, and $(\xi_i)_{i \geq 1}$ are i.i.d. exponentially distributed with rate $\sqrt{2/\pi}$, independent of $(R_i)_{i \geq 1}$. By Wald's identities,

$$\mathbb{E}H = \mathbb{E}N \cdot (\mathbb{E}R_1 - \mathbb{E}\xi_1) = \frac{\pi}{2} \left(\frac{\pi}{2} - \frac{\pi}{2} \right) = 0,$$

and

$$\operatorname{Var} H = \mathbb{E}N \cdot (\operatorname{Var} R_1 + \operatorname{Var} \xi_1) = \frac{\pi}{2} \left(\frac{4 - \pi}{2} + \frac{\pi}{2} \right) = \pi.$$

where $\mathbb{E}R_1$ and $\operatorname{Var} R_1$ are given by (2.2). Moreover,

$$B_{T_1} \stackrel{(d)}{=} -\xi + 1_{\{J < 1\}}H, \quad (4.18)$$

with (ξ, J) independent of H , and the joint distribution of (ξ, J) given by (4.8). So

$$\mathbb{E}B_{T_1} = -\mathbb{E}\xi + \mathbb{P}(J < 1) \cdot \mathbb{E}H = -\sqrt{\frac{\pi}{2}} + \frac{2}{\pi} \cdot 0 = -\sqrt{\frac{\pi}{2}},$$

and

$$\mathbb{E}B_{T_1}^2 = \mathbb{E}\xi^2 - 2\mathbb{E}(\xi 1_{\{J < 1\}}) \cdot \mathbb{E}H + \mathbb{P}(J < 1) \cdot \mathbb{E}H^2 = \pi + 2.$$

□

Remark 4.3. By Leuridan's formula (1.11), the Laplace transform of Δ is given by

$$\Phi_\Delta(\lambda) = -\sum_{n=1}^{\infty} (-\Psi(\lambda))^n \quad \text{for } \lambda > 0 \text{ such that } \Psi(\lambda) < 1, \quad (4.19)$$

where

$$\Psi(\lambda) = \frac{2e^{-\lambda}}{\pi} \int_0^1 \frac{e^{-\lambda t} \sqrt{t}}{t+1} dt + \frac{e^{-2\lambda}}{\pi\lambda}. \quad (4.20)$$

Since $\Psi(\lambda) \rightarrow 0$ as $\lambda \rightarrow \infty$, $\Psi(\lambda) < 1$ for all sufficiently large λ . For such large λ , the expression (4.19) simplifies to

$$\Phi_\Delta(\lambda) = \frac{\Psi(\lambda)}{1 + \Psi(\lambda)}. \quad (4.21)$$

By analytic continuation, the formula (4.21) holds for all $\lambda > 0$. The equality between (4.21) and (1.21) reduces to the following identity

$$\int_0^1 \frac{e^{-\lambda t} \sqrt{t}}{t+1} dt = \frac{\pi e^\lambda}{2} \left[(\operatorname{erf}(\sqrt{\lambda}))^2 - 1 \right] + \sqrt{\frac{\pi}{\lambda}} \operatorname{erf}(\sqrt{\lambda}), \quad (4.22)$$

which does not seem obvious analytically. In an Appendix, we give an analytic proof of (4.22).

To conclude this part, we give another identity in law similar to (1.18).

Proposition 4.4. *Let U be uniform on $[0, 1]$, independent of J and Δ . Then we have the following identity in law*

$$T_1 - 1 \stackrel{(d)}{=} J + 1_{\{J \leq U^2\}} \Delta. \quad (4.23)$$

Proof. Note that $\tilde{T}_1 := T_1 - 1$ is the stationary delay for a renewal process with inter-arrival time distributed according to Δ . If $J = u > 1$ then $\tilde{T}_1 = J$, whereas if $J = u < 1$ then $\tilde{T}_1 = u$ with probability \sqrt{u} , and $\tilde{T}_1 = u + \Delta$ with probability $1 - \sqrt{u}$. This is because a meander of length 1 to the right of time u creates a $(1, 1)$ -minimum for a two-sided Brownian motion at time u if and only if the meander of length u looking backwards from time u to time 0 becomes a meander of length 1 when running further backwards to time $u - 1$. By Brownian excursion theory, the probability that a meander of length u followed by an independent Brownian fragment of length $1 - u$ creates a meander of length 1 is given by

$$\frac{\mathbf{n}(\zeta > 1)}{\mathbf{n}(\zeta > u)} = \frac{\Lambda(1, \infty)}{\Lambda(u, \infty)} = \sqrt{\frac{2}{\pi}} \bigg/ \sqrt{\frac{2}{\pi u}} = \sqrt{u},$$

where $\Lambda(dx)$ is the Lévy measure of a $\frac{1}{2}$ -stable subordinator defined by (2.5). The identity (4.23) follows from the above analysis, where $U \sim \text{Uniform}[0, 1]$ serves as a device to replicate the conditional distribution of \tilde{T}_1 given J . □

By conditioning on J , the identity (4.23) yields a Laplace transform relation, which can be used to provide an alternate derivation of the Laplace transforms of T_1 and of Δ . Though not obviously equivalent, each of the two relations of (1.18) and (4.23) can be derived from the other after substituting in the explicit formula (1.17) for $\Phi_J(\lambda)$, and using the simple density of J on $[0, 1]$. However, neither relation seems to offer much insight into their remarkable implication (1.19).

4.3. A path decomposition between $(1, 1)$ -minima. Let

$$\mathcal{LE}_b := \{t \geq 0 : B_t < B_s \text{ for all } s \in [t, t+b]\},$$

be the set of left ends of forward meanders of length b , and

$$\mathcal{RE}_a := \{t \geq a : B_t < B_s \text{ for all } s \in [t-a, t]\},$$

be the set of right ends of backward meanders of length a . Observe that for $i \geq 1$, $T_i^{a,b} \in \mathcal{LE}_b \cap \mathcal{RE}_a$, and $T_{i+1}^{a,b} = \inf\{t > T_i^{a,b} : t \in \mathcal{LE}_b \cap \mathcal{RE}_a\}$. The following lemma shows that between any $T_i^{a,b}$ and $T_{i+1}^{a,b}$, left ends come before right ends.

Lemma 4.5. *For each $i \geq 1$, let $s \in (T_i^{a,b}, T_{i+1}^{a,b}) \cap \mathcal{LE}_b$ and $t \in (T_i^{a,b}, T_{i+1}^{a,b}) \cap \mathcal{RE}_a$. Then a.s. $s < t$.*

Proof. Suppose by contradiction that there exist $s \in (T_i^{a,b}, T_{i+1}^{a,b}) \cap \mathcal{LE}_b$ and $t \in (T_i^{a,b}, T_{i+1}^{a,b}) \cap \mathcal{RE}_a$ such that $s \geq t$. Let $r := \operatorname{argmin}_{u \in [t,s]} B_u$ be the time at which B attains its a.s. unique minimum between t and s . It is clear that $r \in \mathcal{LE}_b \cap \mathcal{RE}_a$. Thus, $r \geq T_{i+1}^{a,b}$ by definition of $T_{i+1}^{a,b}$. This is impossible since $r \leq s < T_{i+1}^{a,b}$. \square

For $i \geq 1$, let

$$D_i^{a,b} := \inf\{t > T_i^{a,b} : t \in \mathcal{RE}_a\}, \quad (4.24)$$

be the first right end between $T_i^{a,b}$ and $T_{i+1}^{a,b}$, and

$$G_i^{a,b} := \sup\{t < D_i^{a,b} : t \in \mathcal{LE}_b\} \quad (4.25)$$

be the last left end between $T_i^{a,b}$ and $T_{i+1}^{a,b}$. Observe that there are neither left ends nor right ends between $G_i^{a,b}$ and $D_i^{a,b}$. By Lemma 4.5, the next left end after right ends between $T_i^{a,b}$ and $T_{i+1}^{a,b}$ is necessarily a right end; thus is $T_{i+1}^{a,b}$.

Corollary 4.6. *For each $i \geq 1$, $T_{i+1}^{a,b} = \inf\{t > D_i^{a,b} : t \in \mathcal{LE}_b\}$ a.s.*

From now on, we consider the particular case of $a = b = 1$. To simplify notations, write $T_i, \mathcal{LE}, \mathcal{RE}, D_i$ and G_i for $T_i^{1,1}, \mathcal{LE}_1, \mathcal{RE}_1, D_i^{1,1}$ and $G_i^{1,1}$. The following result characterizes the path fragment $B[G_i, D_i]$.

Proposition 4.7. *Almost surely, for each $i \geq 1$,*

- $B_{G_i} = B_{D_i}$ and $D_i - G_i > 1$.
- $B[G_i, D_i] := (B_t - B_{G_i}; G_i \leq t \leq D_i)$ consists of two excursions of lengths smaller than 1.

Proof. Suppose by contradiction that $B_{D_i} < B_{G_i}$. Let $D' := \inf\{t > G_i + 1 : B_t = (B_{D_i} + B_{G_i})/2\}$, and observe that $D' \in \mathcal{RE}$. By path continuity, $D' < D_i$, which contradicts the definition of D_i . Similarly, by considering $G' = \sup\{t < D_i - 1 : B_t = (B_{D_i} + B_{G_i})/2\}$, we exclude the possibility of $B_{D_i} > B_{G_i}$.

Now we argue by contradiction that there exists $u \in (G_i, D_i)$ such that $B_u < B_{G_i} = B_{D_i}$. Let $M := \operatorname{argmin}_{u \in [G_i, D_i]} B_u$ so that $M \in (G_i, D_i)$ and $B_M < B_{G_i} = B_{D_i}$. By definition of G_i , we have $M - G_i > 1$. This implies that $M \in \mathcal{RE}$, which contradicts the definition of D_i . Hence, $B_t \geq B_{G_i} = B_{D_i}$ for all $t \in [G_i, D_i]$, or equivalently $(B_t; G_i \leq t \leq D_i)$ is composed of excursions above the level $B_{G_i} = B_{D_i}$.

If there exists an excursion interval $[u, v] \subset [G_i, D_i]$ with $v - u > 1$, then by path continuity, $[u, v]$ contains at least a left end and a right end. This leads to a contradiction. Further, if there exist $G_i < u < v < D_i$ such that $B_{G_i} = B_u = B_v = B_{D_i}$, then u and v are two local minima at the same level. This is impossible, since a.s. the levels of local minima in Brownian motion are all different, see Kallenberg [27, Lemma 13.15]. Thus, $(B_u - B_{G_i}; G_i \leq t \leq D_i)$ is composed of at most two excursions of lengths no larger than 1.

Finally, observe that for all $t \in (G_i, G_i + 1]$, $B_t > B_{G_i}$ and thereby $t \notin \mathcal{RE}$. This implies that $D_i - G_i \geq 1$. If $D_i - G_i = 1$, then there exists a reflected bridge of length 1 in Brownian motion by a space-time shift. But this is excluded by Pitman and Tang [42, Theorem 4]. \square

According to Proposition 4.7, we get the decomposition (1.29) such that

- $[T_i, T_{i+1}] \cap \mathcal{LE} = [T_i, G_i] \cap \mathcal{LE}$, i.e. left ends of forward meanders of unit length are contained in $[T_i, G_i]$;
- $[T_i, T_{i+1}] \cap \mathcal{RE} = [D_i, T_{i+1}] \cap \mathcal{RE}$, i.e. right ends of backward meanders of unit length are contained in $[D_i, T_{i+1}]$;
- $(G_i, D_i) \cap \mathcal{LE} \cap \mathcal{RE} = \emptyset$, i.e. (G_i, D_i) contains neither left ends of forward meanders nor right ends of backward meanders.

Proof of Theorem 1.6. Observe that T_i is the i^{th} time that the argmin process $(\alpha_t; t \geq 0)$ reaches 0 by a continuous passage from 1. It is obvious that D_i is a stopping time relative to $(\mathcal{G}_t)_{t \geq 0}$, the filtrations of the argmin process α . So $T_{i+1} - D_i$ is independent of $(G_i - T_i, D_i - G_i)$. Further by time reversal of α (Proposition 3.2), we see that $G_i - T_i$, $D_i - G_i$ and $T_{i+1} - D_i$ are mutually independent, and $G_i - T_i \stackrel{(d)}{=} T_{i+1} - D_i$.

By Lemma 3.6, $(1 - \alpha_{D_i+t}; 0 \leq t \leq T_{i+1} - D_i)$ has the same distribution as the age process of excursions above the past-minimum of Brownian motion until the age reaches 1. As seen in Section 2.2, $T_{i+1} - D_i \stackrel{(d)}{=} J$. So $\mathbb{E}(T_{i+1} - D_i) = 1$ by (2.10).

Recall the definitions of $\mu^{\uparrow 1}(x)$, $\Pi^{\uparrow 0}(dx)$ and $s(x, y)$ from (3.7), (3.6) and (3.11). From the Lévy system of the argmin process $(\alpha_t; t \geq 0)$, we have

$$\begin{aligned} \mathbb{P}(D_i - G_i - 1 \in dt)/dt &= c \int_t^1 \Pi^{\uparrow 0}(dx) s(x, x-t) \mu^{\uparrow 1}(x-t) dx \\ &= c \sqrt{\frac{2}{\pi}} \frac{1-t}{(1+t)^2 \sqrt{t}}, \end{aligned}$$

with some constant $c > 0$. Further, $\int_0^1 \mathbb{P}(D_i - G_i - 1 \in dt) = 1$ leads to $c = \sqrt{\pi/2}$. From this follows (1.30). Thus,

$$\mathbb{E}(D_i - G_i) = \int_1^2 t \cdot \frac{2-t}{t^2\sqrt{t-1}} dt = \pi - 2.$$

Alternatively, $\mathbb{E}(D_i - G_i) = \mathbb{E}(T_{i+1} - T_i) - \mathbb{E}(G_i - T_i) - \mathbb{E}(T_{i+1} - D_i) = \pi - 2$ by (1.14), and the fact $\mathbb{E}(G_i - T_i) = \mathbb{E}(T_{i+1} - D_i) = 1$. \square

Remark 4.8. The path decomposition of Theorem 1.6 provides an alternative way to compute the Laplace transform of the inter-arrival time Δ . In fact,

$$\Phi_\Delta(\lambda) = e^{-\lambda} (\Phi_J(\lambda))^2 \int_0^1 \frac{e^{-\lambda t}(1-t)}{(t+1)^2 \sqrt{t}} dt. \quad (4.26)$$

The equality between (4.26) and (1.21) reduces to the following identity

$$\begin{aligned} \int_0^1 \frac{e^{-\lambda t}(1-t)}{\sqrt{t}(1+t)^2} dt &= \pi \lambda e^\lambda \left[(\operatorname{erf}(\sqrt{\lambda})^2 - 1) \right] + 2\sqrt{\pi\lambda} \operatorname{erf}(\sqrt{\lambda}) + e^{-\lambda} \\ &= 2\lambda \int_0^1 \frac{e^{-\lambda t} \sqrt{t}}{t+1} dt + e^{-\lambda}, \end{aligned} \quad (4.27)$$

where the last equality follows from (4.22). The identity (4.27) is easily verified by noting that

$$\begin{aligned} \int_0^1 \frac{e^{-\lambda t}(1-t)}{\sqrt{t}(1+t)^2} dt - 2\lambda \int_0^1 \frac{e^{-\lambda t} \sqrt{t}}{t+1} dt &= \int_0^1 \frac{e^{-\lambda t} [1 - (2\lambda + 1)t - 2\lambda t^2]}{\sqrt{t}(1+t)^2} dt \\ &= \left[\frac{2e^{-\lambda t} \sqrt{t}}{1+t} \right]_0^1 = e^{-\lambda}. \end{aligned}$$

Recall from Section 3.3 that $(l_t^1; t \geq 0)$ is the local times of α at level 1, and $(l_t^0; t \geq 0)$ is the local times of α at level 0. As a consequence of Theorem 1.6, we have

Corollary 4.9.

$$\mathbb{E}l_t^1/t = \mathbb{E}l_t^0/t = \frac{1}{\sqrt{2\pi}} \quad \text{for all } t > 0. \quad (4.28)$$

Proof. By stationarity of the argmin process $(\alpha_t; t \geq 0)$,

$$\mathbb{E}l_t^1/t = \mathbb{E}(l_{[T_i, T_{i+1}]}^1)/\mathbb{E}(T_{i+1} - T_i),$$

where $l_{[T_i, T_{i+1}]}^1$ is the local times of α at level 1 between $[T_i, T_{i+1}]$. Note that $l_{[T_i, T_{i+1}]}^1 = l_{[D_i, T_{i+1}]}^1$. By Lemma 3.6 and Lévy's theorem, $l_{[D_i, T_{i+1}]}^1$ has the same distribution as the first level above which occurs an excursion of length exceeding 1. As seen in Section 2.2, the latter is exponentially distributed with rate $\sqrt{2/\pi}$. Thus, $\mathbb{E}(l_{[T_i, T_{i+1}]}^1) = \sqrt{\pi/2}$. Moreover, $\mathbb{E}(T_{i+1} - T_i) = \pi$ by (1.14). From these follows (4.28). \square

APPENDIX A. ANALYTIC PROOF OF AN IDENTITY

In this appendix, we provide an analytic proof of (4.22). We know from (4.11) that

$$E(e^{-\lambda J} 1_{\{J \leq 1\}}) = \int_0^1 \frac{t^{-1/2}}{\pi} e^{-\lambda t} dt = \frac{\operatorname{erf} \sqrt{\lambda}}{\sqrt{\pi \lambda}},$$

and multiplying this expression by π gives the last term on the right side of (4.22). Since

$$\frac{1}{\sqrt{t}} - \frac{\sqrt{t}}{1+t} = \frac{1}{\sqrt{t}(1+t)},$$

the identity (4.22) is equivalent to the simpler identity

$$\frac{2}{\pi} \int_0^1 \frac{e^{-\lambda t}}{\sqrt{t}(1+t)} dt = e^\lambda \left[1 - (\operatorname{erf} \sqrt{\lambda})^2 \right]. \quad (\text{A.1})$$

The l.h.s of (A.1) is the Laplace transform of a probability density on $(0, 1)$. The term e^λ is just a shift by 1, say $t = v - 1$. So the identity (4.22) is equivalent to the even simpler formula for the Laplace transform of a probability density on $(1, 2)$:

$$\frac{2}{\pi} \int_1^2 \frac{e^{-\lambda v}}{v\sqrt{v-1}} dv = 1 - (\operatorname{erf} \sqrt{\lambda})^2 \quad (\text{A.2})$$

Observe that

$$\int_1^\infty \frac{e^{-\lambda v}}{v\sqrt{v-1}} dv = \pi(1 - \operatorname{erf}(\sqrt{\lambda})). \quad (\text{A.3})$$

By squaring the l.h.s. of (A.3), and making the change of variables $x = v + w$ and $y = v/w$, we get

$$\begin{aligned} \left(\int_1^\infty \frac{e^{-\lambda v}}{v\sqrt{v-1}} dv \right)^2 &= \int_1^\infty \int_1^\infty \frac{e^{-\lambda(v+w)}}{vw\sqrt{v-1}\sqrt{w-1}} dv dw \\ &= \int_2^\infty \frac{e^{-\lambda x}}{x\sqrt{x-1}} \left(\int_{\frac{1}{x-1}}^{x-1} \frac{y+1}{y\sqrt{x-1-y}\sqrt{y-\frac{1}{x-1}}} dy \right) dx \\ &= 2\pi \int_2^\infty \frac{e^{-\lambda x}}{x\sqrt{x-1}} dx. \end{aligned} \quad (\text{A.4})$$

By injecting (A.3) into (A.4), we obtain

$$\int_2^\infty \frac{e^{-\lambda x}}{x\sqrt{x-1}} dx = \frac{\pi}{2} (1 - \operatorname{erf}(\sqrt{\lambda}))^2. \quad (\text{A.5})$$

By subtracting (A.5) from (A.3), and multiplying by $2/\pi$, we get (A.2).

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